

Forchheimer Equations in Porous Media - Part III

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Applied Mathematics Seminar
Texas Tech University
September 15&22, 2010

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- 3 Dependence on the boundary data
- 4 Dependence on the Forchheimer polynomials

Introduction

- Darcy's Law:

$$\alpha u = -\nabla p,$$

- the “two term” law

$$\alpha u + \beta |u| u = -\nabla p,$$

- the “power” law

$$c^n |u|^{n-1} u + a u = -\nabla p,$$

- the “three term” law

$$\mathcal{A}u + \mathcal{B}|u|u + \mathcal{C}|u|^2u = -\nabla p.$$

Here $\alpha, \beta, c, \mathcal{A}, \mathcal{B}$, and \mathcal{C} are empirical positive constants.

General Forchheimer equations

Generalizing the above equations as follows

$$g(|u|)u = -\nabla p.$$

Let $G(s) = sg(s)$. Then $G(|u|) = |\nabla p| \Rightarrow |u| = G^{-1}(|\nabla p|)$. Hence

$$u = -\frac{\nabla p}{g(G^{-1}(|\nabla p|))} = -K(|\nabla p|)\nabla p,$$

where

$$K(\xi) = K_g(\xi) = \frac{1}{g(s)} = \frac{1}{g(G^{-1}(\xi))}, \quad sg(s) = \xi.$$

We derived non-linear Darcy equations from Forchheimer equations.

Equations of Fluids

Let ρ be the density. Continuity equation

$$\frac{d\rho}{dt} = -\nabla \cdot (\rho u),$$

For slightly compressible fluid it takes

$$\frac{d\rho}{dp} = \frac{1}{\kappa} \rho,$$

where $\kappa \gg 1$. Substituting this into the continuity equation yields

$$\frac{d\rho}{dp} \frac{dp}{dt} = -\rho \nabla \cdot u - \frac{d\rho}{dp} u \cdot \nabla p,$$

$$\frac{dp}{dt} = -\kappa \nabla \cdot u - u \cdot \nabla p.$$

Since $\kappa \gg 1$, we neglect the second term in continuity equation

$$\frac{dp}{dt} = -\kappa \nabla \cdot u .$$

Combining the equation of pressure and the Forchheimer equation, one gets after scaling:

$$\frac{dp}{dt} = \nabla \cdot (K(|\nabla p|) \nabla p).$$

Consider the equation on a bounded domain U in \mathbb{R}^3 . The boundary of U consists of two connected components: exterior boundary Γ_2 and interior (accessible) boundary Γ_1 .

- Dirichlet condition on Γ_1 : $p(x, t) = \psi(x, t)$ which is known for $x \in \Gamma_i$.
- On Γ_2 :

$$u \cdot N = 0 \Leftrightarrow \frac{\partial p}{\partial N} = 0.$$

Class $FP(N, \vec{\alpha})$

We introduce a class of “Forchheimer polynomials”

Definition

A function $g(s)$ is said to be of class $FP(N, \vec{\alpha})$ if

$$g(s) = a_0 s^{\alpha_0} + a_1 s^{\alpha_1} + a_2 s^{\alpha_2} + \dots + a_N s^{\alpha_N} = \sum_{j=0}^N a_j s^{\alpha_j},$$

where $N > 0$, $0 = \alpha_0 < \alpha_1 < \alpha_2 < \dots < \alpha_N$, and
 $a_0, a_N > 0$, $a_1, \dots, a_{N-1} \geq 0$. Notation $\alpha_N = \deg(g)$.

Let $N \geq 1$ and $\vec{\alpha}$ be fixed. Let

$$R(N) = \{ \vec{a} = (a_0, a_1, \dots, a_N) \subset \mathbb{R}^{N+1} : a_0, a_N > 0, a_1, \dots, a_{N-1} \geq 0 \}.$$

Let $g = g(s, \vec{a})$ be in $\text{FP}(N, \vec{\alpha})$ with $\vec{a} \in R(N)$. We denote

$$a = \frac{\alpha_N}{\alpha_N + 1} \in (0, 1), \quad b = \frac{a}{2 - a} = \frac{\alpha_N}{\alpha_N + 2} \in (0, 1),$$

$$\chi(\vec{a}) = \max \left\{ a_0, a_1, \dots, a_N, \frac{1}{a_0}, \frac{1}{a_N} \right\} \geq 1.$$

Many estimates below will have constants depending on $\chi(\vec{a})$.

Lemma

Let $g(s, \vec{a})$ be in class $FP(N, \vec{\alpha})$. One has for any $\xi \geq 0$ that

$$\frac{C_0^{-1} \chi(\vec{a})^{-1-a}}{(1+\xi)^a} \leq K(\xi, \vec{a}) \leq \frac{C_0 \chi(\vec{a})^{1+a}}{(1+\xi)^a},$$

and for any $m \geq 1, \delta > 0$ that

$$C_0^{-1} \chi(\vec{a})^{-1-a} \frac{\delta^a}{(1+\delta)^a} (\xi^{m-a} - \delta^{m-a}) \leq K(\xi, \vec{a}) \xi^m \leq C_0 \chi(\vec{a})^{1+a} \xi^{m-a},$$

where $C_0 = C_0(N, \alpha_N)$ depends on N, α_N only.

In particular, when $m = 2, \delta = 1$, one has

$$2^{-a} C_0^{-1} \chi(\vec{a})^{-1-a} (\xi^{2-a} - 1) \leq K(\xi, \vec{a}) \xi^2 \leq C_0 \chi(\vec{a})^{1+a} \xi^{2-a}.$$

The Monotonicity

Proposition

(i) For any $y, y' \in \mathbb{R}^n$, one has

$$(K(|y|, \vec{a})y - K(|y'|, \vec{a})y') \cdot (y - y') \geq aK(|y| \vee |y'|, \vec{a})|y - y'|^2.$$

(ii) For any functions p_1 and p_2 one has

$$\begin{aligned} & \int_U (K(|\nabla p_1|, \vec{a})\nabla p_1 - K(|\nabla p_2|, \vec{a})\nabla p_2) \cdot (\nabla p_1 - \nabla p_2) dx \\ & \geq a \left(\int_U K(|\nabla p_1| \vee |\nabla p_2|, \vec{a}) |\nabla p_1 - \nabla p_2|^2 dx \right) \\ & \geq C_5 \left(\int_U |\nabla p_1 - \nabla p_2|^{2-a} dx \right)^{\frac{2}{2-a}} (1 + \|\nabla p_1\|_{L^{2-a}(U)} \vee \|\nabla p_2\|_{L^{2-a}(U)})^{-a}, \end{aligned}$$

where $C_5 = C_5(N, \deg(g), \chi(\vec{a}))$.

Poincaré–Sobolev inequality

Let $f(x)$ and $\xi(x)$ be two functions on U with $f(x)$ vanishing on Γ_1 and $\xi(x) \geq 0$. Then

$$\left(\int_U |f(x)|^{(2-a)^*} dx \right)^{\frac{2}{(2-a)^*}} \leq C_6 \left(\int_U K(\xi(x), \vec{a}) |\nabla f(x)|^2 dx \right) \\ \times \left(1 + \int_U H(\xi(x), \vec{a}) dx \right)^{\frac{a}{2-a}},$$

where $(2-a)^* = n(2-a)/(n-(2-a))$ and $C_6 = C_6(N, \deg(g), \chi(\vec{a}), U)$. Subsequently, when $\deg(g) \leq 4/(n-2)$ one has

$$\int_U |f(x)|^2 dx \leq C_7 \left(\int_U K(\xi(x), \vec{a}) |\nabla f(x)|^2 dx \right) \left(1 + \int_U H(\xi(x), \vec{a}) dx \right)^{\frac{a}{2-a}},$$

where $C_7 = C_7(N, \deg(g), \chi(\vec{a}), U)$.

Degree Condition (DC)

$$\deg(g) \leq \frac{4}{n-2} \iff 2 \leq (2-a)^* = \frac{n(2-a)}{n-(2-a)}.$$

Extension of the boundary data

Let $\Psi(x, t)$, $x \in U$, $t \in [0, \infty)$ be an extension of $\psi(x, t)$ from Γ_1 to U . For instance, one can use the following harmonic extension Ψ of ψ :

$$\Delta \Psi = 0 \quad \text{on } U, \quad \Psi \Big|_{\Gamma_1} = \psi \quad \text{and} \quad \Psi \Big|_{\Gamma_2} = 0.$$

We denote such Ψ by $\mathcal{H}(\psi)$. Then the Sobolev norms of $\mathcal{H}(\psi)$ on U can be bounded by the Sobolev or Besov norms of ψ on Γ_1 . In particular, we have

$$\left\| \frac{\partial^k}{\partial t^k} \mathcal{H}(\psi) \right\|_{W^{r,2}(U)} \leq C(k, r) \left\| \frac{\partial^k}{\partial t^k} \psi \right\|_{W^{r,2}(\Gamma_1)},$$

for $k = 0, 1, 2$, $r = 0, 1$.

Shift of solutions

Shifted solution: Let $\bar{p} = p - \Psi$, then \bar{p} satisfies

$$\frac{\partial \bar{p}}{\partial t} = \nabla \cdot (K(|\nabla p|) \nabla p) - \Psi_t \quad \text{on } U \times (0, \infty),$$
$$\bar{p} = 0 \quad \text{on } \Gamma_1 \times (0, \infty).$$

Define:

$$H(\xi, \vec{a}) = \int_0^{\xi^2} K(\sqrt{s}, \vec{a}) ds.$$

We denote $H(x, t) = H[p](x, t) = H(|\nabla p(x, t)|, \vec{a})$.
Constants in estimates below depend on $\chi(\vec{a})$.

A priori Estimates - I(a)

Lemma

One has

$$\frac{1}{2} \frac{d}{dt} \int_U \bar{p}^2(x, t) dx \leq -C \int_U H(x, t) dx + CG_1(t),$$

where

$$G_1(t) = \int_U |\nabla \Psi(x, t)|^2 dx + \left(\int_U |\Psi_t(x, t)|^{r_0} dx \right)^{\frac{2-a}{r_0(1-a)}} + \left(\int_U |\Psi_t(x, t)|^{r_0} dx \right)^{\frac{1}{r_0}}$$

with r_0 denoting the conjugate exponent of $(2-a)^*$, thus explicitly having the value

$$r_0 = \frac{n(2-a)}{(2-a)(n+1)-n} = \frac{n(2+\alpha_N)}{n+2+\alpha_N}.$$

A priori Estimates - I(b)

Corollary

One has for $t \geq 0$ that

$$\int_U \bar{p}^2(x, t) dx \leq \int_U \bar{p}^2(x, 0) dx + C\Lambda_1(t),$$

where

$$\Lambda_1(t) = \int_0^t G_1(\tau) d\tau.$$

In the case $\deg(g) \leq 4/(n-2)$ one has

$$\int_U \bar{p}^2(x, t) dx \leq \int_U \bar{p}^2(x, 0) dx + C\Lambda_2(t),$$

where $\Lambda_2(t) = (1 + Env(G_1)(t))^{\frac{2}{2-a}}$, with $Env(G_1)(t)$ being a continuous, increasing envelop of the function $G_1(t)$ on $[0, \infty)$.

A priori Estimates - I: Proofs

- No Degree Condition: minimal information
- With Degree Condition: Sobolev-Poincare inequality and non-linear differential inequality (weaker than the usual Gronwall's)

Non-linear differential inequality

Lemma

Suppose

$$y' \leq -Ay^\alpha + f(t),$$

for all $t > 0$, with $A, \alpha > 0$ and $y(t), f(t) \geq 0$.

Let $F(t)$ be a continuous, increasing envelop of $f(t)$ on $[0, \infty)$. Then one has

$$y(t) \leq y(0) + A^{-1/\alpha} F(t)^{1/\alpha}, \quad \forall t \geq 0.$$

A priori Estimates - II(a)

Lemma

For any $\varepsilon > 0$, one has

$$\frac{d}{dt} \int_U H(x, t) dx + \int_U \bar{p}_t^2(x, t) dx \leq \varepsilon \int_U H(x, t) dx + C_\varepsilon G_2(t),$$

where C_ε is positive and

$$G_2(t) = \int_U |\nabla \Psi_t(x, t)|^2 dx + \int_U |\Psi_t(x, t)|^2 dx.$$

Consequently, one has

$$\frac{d}{dt} \left[\int_U H(x, t) + \bar{p}_t^2(x, t) dx \right] + \int_U \bar{p}_t^2(x, t) dx \leq -C \int_U H(x, t) dx + CG_3(t),$$

where $G_3(t) = G_1(t) + G_2(t)$.

A priori Estimates - II(b)

Corollary

(i) Given $\delta > 0$, there is $C_\delta > 0$ such that for all $t \geq 0$ one has

$$\int_U H(x, t) dx \leq e^{\delta t} \int_U H(x, 0) dx + C_\delta \int_0^t e^{\delta(t-\tau)} G_2(\tau) d\tau.$$

(ii) For $t \geq 0$, one has

$$\begin{aligned} & \int_U H(x, t) + \bar{p}^2(x, t) dx + \boxed{\int_0^t \int_U \bar{p}_t^2(x, \tau) dx d\tau} \\ & \leq \int_U H(x, 0) + \bar{p}^2(x, 0) dx + C \int_0^t G_3(\tau) d\tau. \end{aligned}$$

A priori Estimates - II(c)

Corollary

For $t \geq 0$, one has

$$\begin{aligned}\int_U H(x, t) dx &\leq e^{-C_1 t} \int_U H(x, 0) dx + C \int_U \bar{p}^2(x, 0) dx \\ &\quad + C \int_0^t e^{-C_1(t-\tau)} (\Lambda_1(\tau) + G_3(\tau)) d\tau.\end{aligned}$$

In case $\deg(g) \leq 4/(n-2)$, one has for $t \geq 0$ that

$$\begin{aligned}\int_U H(x, t) dx &\leq e^{-C_1 t} \int_U H(x, 0) dx + C \int_U \bar{p}^2(x, 0) dx \\ &\quad + C \int_0^t e^{-C_1(t-\tau)} (\Lambda_2(\tau) + G_3(\tau)) d\tau.\end{aligned}$$

A priori Estimates - II(d)

Lemma

Suppose $\deg(g) \leq \frac{4}{n-2}$. Then

$$\int_U \bar{p}^2(x, t) dx \leq Ch(t) \left(\int_U H(x, t) dx + \int_U |\nabla \Psi(x, t)|^2 dx \right),$$

$$\left(\int_U \bar{p}^2(x, t) dx \right)^{\frac{2-a}{2}} \leq C \left(1 + \int_U H(x, t) dx \right) + C \int_U |\nabla \Psi(x, t)|^2 dx,$$

where

$$h(t) = \left(1 + \int_U H(x, t) dx \right)^{\frac{a}{2-a}}.$$

A priori Estimates - II(e)

Proposition

Suppose $\deg(g) \leq \frac{4}{n-2}$. One has the following two estimates

$$\begin{aligned} \int_U H(x, t) + \bar{p}^2(x, t) dx &\leq e^{-C_1 \int_0^t h^{-1}(\tau) d\tau} \left(\int_U H(x, 0) + \bar{p}^2(x, 0) dx \right) \\ &\quad + C \int_0^t e^{-C_1 \int_\tau^t h^{-1}(\theta) d\theta} G_3(\tau) d\tau, \end{aligned}$$

and

$$\int_U H(x, t) + \bar{p}^2(x, t) dx \leq \int_U H(x, 0) + \bar{p}^2(x, 0) dx + C (1 + Env(G_3)(t))^{\frac{2}{2-a}}.$$

A priori Estimates - III(a)

Let $q = p_t$, $\bar{q} = \bar{p}_t$.

Lemma

One has for $t > 0$ that

$$\frac{d}{dt} \int_U \bar{q}^2 dx \leq -C \int_U K(|\nabla p|) |\nabla \bar{q}|^2 dx + C \int_U |\nabla \Psi_t|^2 dx + \int_U |\bar{q} \Psi_{tt}| dx.$$

Note: May happen

$$\limsup_{t \rightarrow 0} \int_U \bar{p}_t^2(x, t) dx = \infty.$$

A priori Estimates - III(b)

Proposition

One has for $t \geq t_0 > 0$ that

$$\begin{aligned} \int_U \bar{p}_t^2(x, t) dx &\leq t_0^{-1} \left[\int_U H(x, 0) + \bar{p}^2(x, 0) dx + C \int_0^{t_0} G_3(\tau) d\tau \right] \\ &\quad + C \int_0^t \left\{ \int_U |\nabla \Psi_t(x, \tau)|^2 dx + h(\tau) \left(\int_U |\Psi_{tt}(x, \tau)|^{r_0} dx \right)^{\frac{2}{r_0}} \right\} d\tau. \end{aligned}$$

If $\deg(g) \leq \frac{4}{n-2}$, then for $t \geq t_0 > 0$ one has

$$\begin{aligned} \int_U \bar{p}_t^2(x, t) dx &\leq t_0^{-1} e^{-C_1 \int_{t_0}^t \frac{1}{h(\tau)} d\tau} \left[\int_U H(x, 0) + \bar{p}^2(x, 0) dx + C \int_0^{t_0} G_3(\tau) d\tau \right] \\ &\quad + C \int_0^t e^{-C_1 \int_\tau^t \frac{1}{h(\theta)} d\theta} \left(\int_U |\nabla \Psi_t(x, \tau)|^2 dx + h(\tau) \int_U |\Psi_{tt}(x, \tau)|^2 dx \right) d\tau. \end{aligned}$$

A priori Estimates - IV(a)

Proposition

Suppose $\deg(g) \leq 4/(n-2)$.

(i) For any $t \geq t_0 > 0$ one has

$$\begin{aligned} \int_U H(x, t) + \bar{p}_t^2(x, t) + \bar{p}^2(x, t) dx &\leq (1 + t_0^{-1}) e^{-C_1 \int_{t_0}^t h^{-1}(\tau) d\tau} \left[\int_U H(x, 0) \right. \\ &\quad \left. + \bar{p}^2(x, 0) dx + C \int_0^{t_0} G_3(\tau) d\tau \right] + C \int_0^t e^{-C_1 \int_\tau^t h^{-1}(\theta) d\theta} G_4(\tau) d\tau, \end{aligned}$$

where $G_4(t) = G_3(t) + \int_U \Psi_{tt}^2(x, t) dx$.

A priori Estimates - IV(b)

Proposition (continued)

(ii) Assume, in addition, that

$$M_0 \stackrel{\text{def}}{=} \int_0^\infty e^{-C_1(t-\tau)} (\Lambda_2(t) + G_3(t)) d\tau < \infty.$$

Then there is $d_0 > 0$ depending on the initial data of the solution and the value M_0 above so that

$$\begin{aligned} \int_U H(x, t) + \bar{p}_t^2(x, t) + \bar{p}^2(x, t) dx &\leq (1 + t_0^{-1}) e^{-d_0(t-t_0)} \left[\int_U H(x, 0) \right. \\ &\quad \left. + \bar{p}^2(x, 0) dx + C \int_0^{t_0} G_3(\tau) d\tau \right] + C \int_0^t e^{-d_0(t-\tau)} G_4(\tau) d\tau, \end{aligned}$$

for all $t \geq t_0 > 0$.

A priori Estimates - IV(c): No condition on $h(t)$

Proposition

Suppose $\deg(g) \leq \frac{4}{n-2}$. One has for $t \geq t_0 > 0$ that

$$\begin{aligned} & \int_U H(x, t) + \bar{p}_t^2(x, t) + \bar{p}^2(x, t) dx \\ & \leq (1 + t_0^{-1}) \int_U H(x, 0) + \bar{p}^2(x, 0) dx + C(1 + Env(G_4)(t))^{\frac{2}{2-a}}. \end{aligned}$$

A priori Estimates - V(a): Uniform bounds

Theorem

Suppose $\deg(g) \leq 4/(n-2)$. Assume that

$$\sup_{[0,\infty)} \|\psi(\cdot, t)\|_{W^{1,2}(\Gamma_1)}, \quad \sup_{[0,\infty)} \|\psi_t(\cdot, t)\|_{W^{1,2}(\Gamma_1)} < \infty.$$

Then

$$\begin{aligned} \sup_{[0,\infty)} \int_U H(x, t) + p^2(x, t) dx &\leq 4 \int_U H(x, 0) + p^2(x, 0) dx + C \left(1 + \right. \\ &\quad \left. \sup_{[0,\infty)} \|\psi(\cdot, t)\|_{W^{1,2}(\Gamma_1)}^2 + \sup_{[0,\infty)} \|\psi_t(\cdot, t)\|_{L^2(\Gamma_1)}^{\frac{2-a}{1-a}} + \sup_{[0,\infty)} \|\psi_t(\cdot, t)\|_{W^{1,2}(\Gamma_1)}^2 \right)^{\frac{2}{2-a}}. \end{aligned}$$

A priori Estimates - V(b): Uniform bounds

Theorem (continued)

If, in addition,

$$\sup_{[0,\infty)} \|\psi_{tt}(\cdot, t)\|_{L^2(\Gamma_1)} < \infty,$$

then for any $t_0 > 0$ one has

$$\begin{aligned} & \sup_{[t_0,\infty)} \int_U H(x, t) + p_t^2(x, t) + p^2(x, t) dx \leq C(1 + t_0^{-1}) \int_U H(x, 0) \\ & + p^2(x, 0) dx + Ct_0^{-1} \int_{\Gamma_1} |\psi(x, 0)|^2 d\sigma + C \left(1 + \sup_{[0,\infty)} \|\psi(\cdot, t)\|_{W^{1,2}(\Gamma_1)}^2 \right. \\ & \left. + \sup_{[0,\infty)} \|\psi_t(\cdot, t)\|_{L^2(\Gamma_1)}^{\frac{2-a}{1-a}} + \sup_{[0,\infty)} \|\psi_t(\cdot, t)\|_{W^{1,2}(\Gamma_1)}^2 + \sup_{[0,\infty)} \|\psi_{tt}(\cdot, t)\|_{L^2(\Gamma_1)}^2 \right)^{\frac{2}{2-a}}. \end{aligned}$$

Dependence on the boundary data

Let $p_1(x, t)$ and $p_2(x, t)$ be two solutions of the IBVP with the boundary profiles $\psi_1(x, t)$ and $\psi_2(x, t)$, respectively. Let $\Psi_k(x, t)$ be an extension of $\psi_k(x, t)$, for $k = 1, 2$.

We denote

$$z(x, t) = p_1(x, t) - p_2(x, t), \quad \Psi(x, t) = \Psi_1(x, t) - \Psi_2(x, t),$$

$$\bar{p}_k = p_k - \Psi_k, \quad k = 1, 2, \quad \bar{z} = \bar{p}_1 - \bar{p}_2 = z - \Psi.$$

Let $H_k(x, t) = H[p_k](x, t) = H(|\nabla p_k(x, t)|, \vec{a})$ for $k = 1, 2$.

Recall that $a = \frac{\deg(g)}{\deg(g)+1}$ and $b = \frac{a}{2-a} = \frac{\deg(g)}{\deg(g)+2}$.

We will establish various estimates for $\bar{Z}(t) \stackrel{\text{def}}{=} \int_U \bar{z}^2(x, t) dx$, for $t \geq 0$.

First, we derive a general differential inequality for $\bar{Z}(t)$.

Lemma

One has for all $t > 0$,

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_U \bar{z}^2 dx &\leq -C \left(\int_U |\nabla \bar{z}|^{2-a} dx \right)^{\frac{2}{2-a}} (1 + \|H_1\|_{L^1} + \|H_2\|_{L^1})^{-b} \\ &\quad + C(1 + \|H_1\|_{L^1} + \|H_2\|_{L^1})^{-b} \|\nabla \Psi\|_{L^{2-a}}^2 \\ &\quad + C(\|H_1\|_{L^1} + \|H_2\|_{L^1})^{1/2} \|\nabla \Psi\|_{L^2} \\ &\quad + C(1 + \|H_1\|_{L^1} + \|H_2\|_{L^1})^b \|\Psi_t\|_{L^{r_0}}^2, \end{aligned}$$

where r_0 is the conjugate of $(2-a)^$.*

Let $G_j[\Psi_k]$, $k = 1, 2$, $j = 1, 2, 3, 4$, denote the quantity G_j for corresponding solution p_k with boundary data extension Ψ_k . Similarly, let $\Lambda_j[\Psi_k]$, $k = 1, 2$, $j = 1, 2$, denote the corresponding quantity Λ_j defined for Ψ_k .

Let $\bar{m}(t) = \bar{m}_1(t) + \bar{m}_2(t)$, where for $k = 1, 2$,

$$\begin{aligned}\bar{m}_k(t) &= e^{-C_1 t} \int_U H_k(x, 0) dx + \int_U \bar{p}_k^2(x, 0) dx \\ &\quad + \int_0^t e^{-C_1(t-\tau)} \left(\Lambda_1[\Psi_k](\tau) + G_3[\Psi_k](\tau) \right) d\tau.\end{aligned}$$

Estimate for finite time interval

Theorem

One has for all $t \geq 0$ that

$$\begin{aligned} \int_U \bar{z}^2(x, t) dx &\leq \int_U \bar{z}^2(x, 0) dx + C \int_0^t \left(\|\nabla \Psi(\cdot, \tau)\|_{L^{2-a}}^2 \right. \\ &\quad \left. + \bar{m}(\tau)^{1/2} \|\nabla \Psi(\cdot, \tau)\|_{L^2} + (1 + \bar{m}(\tau))^b \|\Psi_t(\cdot, \tau)\|_{L^{r_0}}^2 \right) d\tau. \end{aligned}$$

Consequently, for any given $T > 0$,

$$\begin{aligned} \sup_{[0, T]} \int_U z^2(x, t) dx &\leq 4 \int_U z^2(x, 0) dx + 6 \sup_{[0, T]} \|\Psi(\cdot, t)\|_{L^2}^2 + CT \sup_{[0, T]} \|\nabla \Psi(\cdot, t)\|_{L^{2-a}}^2 \\ &\quad + CT(1 + A_* + D_*(T))^\delta \left(\sup_{[0, T]} \|\nabla \Psi(\cdot, t)\|_{L^2} + \sup_{[0, T]} \|\Psi_t(\cdot, t)\|_{L^{r_0}}^2 \right), \end{aligned}$$

Above, $\delta = \max\{1/2, b\}$,

$$A_* = \int_U H_1(x, 0) dx + \int_U \bar{p}_1^2(x, 0) dx + \int_U H_2(x, 0) dx + \int_U \bar{p}_2^2(x, 0) dx,$$

$$D_*(T) = \sum_{k=1}^2 \sup_{[0, T]} \int_0^t e^{-C_1(t-\tau)} \left(\Lambda_1[\Psi_k](\tau) + G_3[\Psi_k](\tau) \right) d\tau.$$

Corollary

Suppose for both $k = 1, 2$ and $t \geq 0$ one has

$$\|\nabla \Psi_k(\cdot, t)\|_{L^2}^2 + \|(\Psi_k)_t(\cdot, t)\|_{L^{r_0}}^{\frac{2-a}{1-a}} \leq C(1+t)^{r_1}$$

and

$$\|\nabla(\Psi_k)_t(\cdot, t)\|_{L^2}^2 + \|(\Psi_k)_t(\cdot, t)\|_{L^2}^2 \leq C(1+t)^{r_2},$$

where $r_1, r_2 > 0$. Let $r_3 = 1 + \max\{r_1 + 1, r_2\}$. Then

$$\begin{aligned} \int_U \bar{z}^2(x, t) dx &\leq \int_U \bar{z}^2(x, 0) dx \\ &+ C(1+A_*) \int_0^t (1+\tau)^{r_3/2} \|\nabla \Psi(\cdot, \tau)\|_{L^2} + (1+\tau)^{r_3 b} \|\Psi_t(\cdot, \tau)\|_{L^r}^2 d\tau, \end{aligned}$$

where A_* is defined previously.

Estimate for all time under the Degree Condition

Using appropriate estimates of $\int_U H_k(x, t)dx$, we denote

$$\begin{aligned} m_k(t) &= e^{-C_1 t} \int_U H_k(x, 0)dx + \int_U \bar{p}_k^2(x, 0)dx \\ &+ \int_0^t e^{-C_1(t-\tau)} \left(\Lambda_2[\Psi_k](\tau) + G_3[\Psi_k](\tau) \right) d\tau, \quad k = 1, 2. \end{aligned}$$

Also, let

$$m(t) = m_1(t) + m_2(t) \text{ and } S(t', t) = \int_{t'}^t (1 + m(\tau))^{-b} d\tau.$$

Theorem

Suppose $\deg(g) \leq 4/(n - 2)$. Then for all $t \geq 0$ one has

$$\begin{aligned} \int_U \bar{z}^2(x, t) dx &\leq e^{-C_1 S(0, t)} \int_U \bar{z}^2(x, 0) dx \\ &+ C \int_0^t e^{-C_1 S(\tau, t)} \left(\|\nabla \Psi(\cdot, \tau)\|_{L^{2-a}}^2 \right. \\ &\quad \left. + m(\tau)^{1/2} \|\nabla \Psi(\cdot, \tau)\|_{L^2} + (1 + m(\tau))^b \|\Psi_t(\cdot, \tau)\|_{L^{r_0}}^2 \right) d\tau. \end{aligned}$$

Corollary

Suppose $\deg(g) \leq 4/(n - 2)$. Assume that

$$\sum_{k=1}^2 \left(\sup_{[0,\infty)} \Lambda_2[\Psi_k](t) + \sup_{[0,\infty)} G_3[\Psi_k](t) \right) < \infty.$$

Then one has

$$\begin{aligned} \sup_{[0,\infty)} \int_U z^2(x, t) dx &\leq 4 \int_U z^2(x, 0) dx + C \sup_{[0,\infty)} \|\Psi(\cdot, t)\|_{L^2}^2 \\ &+ C \sup_{[0,\infty)} \left(A_*^b \|\nabla \Psi(\cdot, t)\|_{L^{2-a}}^2 + A_*^{b+\frac{1}{2}} \|\nabla \Psi(\cdot, t)\|_{L^2} + A_*^{2b} \|\Psi_t(\cdot, t)\|_{L^{r_0}}^2 \right), \end{aligned}$$

where

$$A_* = 1 + \sum_{k=1}^2 \left(\int_U |\nabla p_k(x, 0)|^{2-a} + \bar{p}_k^2(x, 0) dx + \sup_{[0,\infty)} \Lambda_2[\Psi_k](t) + \sup_{[0,\infty)} G_3[\Psi_k](t) \right)$$

Corollary

Suppose $\deg(g) \leq 4/(n-2)$. Assume that

$$\lim_{t \rightarrow \infty} S(0, t) = \infty,$$

$$\lim_{t \rightarrow \infty} (1 + m(t))^{b+\frac{1}{2}} \|\nabla \Psi(\cdot, t)\|_{L^2} = 0, \quad \lim_{t \rightarrow \infty} (1 + m(t))^b \|\Psi_t(\cdot, t)\|_{L^0} = 0.$$

Then $\lim_{t \rightarrow \infty} \int_U \bar{z}^2(x, t) dx = 0$.

Dependence on the Forchheimer polynomials

Let the boundary data $\psi(x, t)$ on Γ_1 be fixed. For each coefficient vector \vec{a} , we denote by $p(x, t; \vec{a})$ the solution of

$$p_t = \nabla \cdot (K(|\nabla p|) \nabla p), \quad p|_{\Gamma_1} = \psi, \quad \frac{\partial p}{\partial N} \Big|_{\Gamma_2} = 0,$$

with $K = K(\xi, \vec{a})$.

Let $g_1 = g(s, \vec{a}^{(1)})$ and $g_2 = g(s, \vec{a}^{(2)})$ be two functions of class $\text{FP}(N, \vec{\alpha})$.

Let $p_k = p_k(x, t; \vec{a}^{(k)})$ for $k = 1, 2$. Let $p = p_1 - p_2$, then

$$\frac{\partial p}{\partial t} = \nabla \cdot (K(|\nabla p_1|, \vec{a}^{(1)}) \nabla p_1) - \nabla \cdot (K(|\nabla p_2|, \vec{a}^{(2)}) \nabla p_2).$$

Multiplying this equation by p and integrating by parts over the domain yield

$$\frac{1}{2} \frac{d}{dt} \int_U p^2 dx = - \int_U (K(|\nabla p_1|, \vec{a}^{(1)}) \nabla p_1 - K(|\nabla p_2|, \vec{a}^{(2)}) \nabla p_2) \cdot (\nabla p_1 - \nabla p_2) dx$$

Perturbed Monotonicity

Let \vec{a} and \vec{a}' be two arbitrary vectors. We denote by $\vec{a} \vee \vec{a}'$ and $\vec{a} \wedge \vec{a}'$ the maximum and minimum vectors of the two, respectively, with components

$$(\vec{a} \vee \vec{a}')_j = \max\{a_j, a'_j\} \quad \text{and} \quad (\vec{a} \wedge \vec{a}')_j = \min\{a_j, a'_j\}.$$

Then component-wise one has $\vec{a} \wedge \vec{a}' \leq \vec{a}, \vec{a}' \leq \vec{a} \vee \vec{a}'$.

Define $\chi(\vec{a}, \vec{a}') = \max\{\chi(\vec{a}), \chi(\vec{a}')\}$. Note that

$$\chi(\vec{a} \vee \vec{a}'), \quad \chi(\vec{a} \wedge \vec{a}') \leq \chi(\vec{a}, \vec{a}'),$$

$$\chi(t\vec{a} + (1-t)\vec{a}') \leq \chi(\vec{a}, \vec{a}') \quad \forall t \in [0, 1].$$

Perturbed Monotonicity

Lemma

Let $g(s, \vec{a})$ and $g(s, \vec{a}')$ belong to class $FP(N, \vec{\alpha})$. Then for any y, y' in \mathbb{R}^n , one has

$$(K(|y|, \vec{a})y - K(|y'|, \vec{a}')y') \cdot (y - y') \geq (1 - a)K(|y| \vee |y'|, \vec{a} \vee \vec{a}')|y - y'|^2 - N\chi(\vec{a}, \vec{a}')|\vec{a} - \vec{a}'|K(|y| \vee |y'|, \vec{a} \wedge \vec{a}')(|y| \vee |y'|)|y - y'|,$$

where $a \in (0, 1)$.

Let

$$\begin{aligned}\overline{M}_k(t) = & 1 + e^{-C_1 t} \int_U |\nabla p_k(x, 0)|^{2-a} dx + \int_U \bar{p}_k^2(x, 0) dx \\ & + \int_0^t e^{-C_1(t-\tau)} (\Lambda_1(\tau) + G_3(\tau)) d\tau, \quad k = 1, 2,\end{aligned}$$

and $\overline{M} = \overline{M}_1 + \overline{M}_2$. One has

$$\int_U K(|\nabla p_k|, \vec{a}^{(1)} \wedge \vec{a}^{(2)}) |\nabla p_k|^2 dx \leq C + C \int_U |\nabla p_k|^{2-a} dx \leq C \overline{M}_k,$$

where C depends on $\chi(\vec{a}^{(1)} \wedge \vec{a}^{(2)})$ and $\chi(\vec{a}^{(k)})$.

Denote $H_k(x, t) = H(|\nabla p_k(x, t)|, \vec{a}^{(k)})$ for $k = 1, 2$.

Finite time estimate

First, we obtain the continuous dependence for finite time.

Proposition

For $t \geq 0$ one has

$$\begin{aligned} \int_U |p_1(x, t) - p_2(x, t)|^2 dx &\leq \int_U |p_1(x, 0) - p_2(x, 0)|^2 dx \\ &+ C |\vec{a}^{(1)} - \vec{a}^{(2)}| \int_0^t \overline{M}(\tau) d\tau, \end{aligned}$$

where $C > 0$ depends on N , α_N , and $\chi(\vec{a}^{(1)}, \vec{a}^{(2)})$. Consequently, the solution $p(x, t; \vec{a})$ depends continuously (in finite time intervals) on the initial data and the coefficient vector $\vec{a} \in R(N)$.

Large time estimate under the Degree Condition

Under the Degree Condition, $\int_U |\nabla p_k(x, t)|^{2-a} dx$ is bounded by $CM_k(t)$ where

$$\begin{aligned} M_k(t) = & 1 + e^{-C_1 t} \int_U |\nabla p_k(x, 0)|^{2-a} dx + \int_U \bar{p}_k^2(x, 0) dx \\ & + \int_0^t e^{-C_1(t-\tau)} (\Lambda_2(\tau) + G_3(\tau)) d\tau, \quad k = 1, 2, \end{aligned}$$

and $M(t) = M_1(t) + M_2(t)$.

Proposition

Suppose $\alpha_N \leq 4/(n - 2)$. Then for $t \geq 0$,

$$\int_U |p_1(x, t) - p_2(x, t)|^2 dx \leq e^{-C_1 \int_0^t M(\tau)^{-b}(\tau) d\tau} \int_U |p_1(x, 0) - p_2(x, 0)|^2 dx \\ + C_2 |\vec{a}^{(1)} - \vec{a}^{(2)}| \int_0^t e^{-C_1 \int_\tau^t M(\theta)^{-b}(\theta) d\theta} M(\tau) d\tau.$$

Assume, in addition, that $M_0 \stackrel{\text{def}}{=} \sup_{[0, \infty)} M(t) < \infty$ then

$$\int_U |p_1(x, t) - p_2(x, t)|^2 dx \leq e^{-C_1 t / M_0^b} \int_U |p_1(x, 0) - p_2(x, 0)|^2 dx \\ + C_2 C_1 M_0^{1+b} |\vec{a}^{(1)} - \vec{a}^{(2)}|$$

for all $t \geq 0$, and consequently

$$\limsup_{t \rightarrow \infty} \int_U |p_1(x, t) - p_2(x, t)|^2 dx \leq C_2 C_1 M_0^{1+b} |\vec{a}^{(1)} - \vec{a}^{(2)}|.$$

Uniform estimate in the coefficients

Let D be a compact set in $R(N)$. Define

$$\hat{\chi}(D) = \max\{\chi(\vec{a}), \vec{a} \in D\}.$$

Note that $\hat{\chi}(D) \in (0, \infty)$ and for any $\vec{a} \in D$, one has

$$\hat{\chi}(D)^{-1} \leq \chi(\vec{a})^{-1} \leq \chi(\vec{a}) \leq \hat{\chi}(D).$$

Let $\vec{a}^{(k)}$ belong to D for $k = 1, 2$. Set

$$A_* = 2 + \sum_{k=1}^2 \left(\int_U |\nabla p_k(x, 0)|^{2-a} dx + \int_U \bar{p}_k^2(x, 0) dx \right)$$

$$\bar{\lambda}(t) = 1 + \int_0^t e^{-C_1(t-\tau)} (\Lambda_1(\tau) + G_3(\tau)) d\tau,$$

$$\lambda(t) = 1 + \int_0^t e^{-C_1(t-\tau)} (\Lambda_2(\tau) + G_3(\tau)) d\tau,$$

where C_1 depends on N , α_N and $\hat{\chi}(D)$.

Then $\bar{M}(t) \leq A_* \bar{\lambda}(t)$, $M(t) \leq A_* \lambda(t)$.

Theorem

Assume, additionally, that $\alpha_N \leq 4/(n-2)$.

(i) One has for $t \geq 0$ that

$$\begin{aligned} \int_U |p_1(x, t) - p_2(x, t)|^2 dx &\leq e^{-C_1 A_*^{-b} \int_0^t \lambda(\tau)^{-b} d\tau} \int_U |p_1(x, 0) - p_2(x, 0)|^2 dx \\ &+ C_2 A_* |\vec{a}^{(1)} - \vec{a}^{(2)}| \int_0^t e^{-C_1 A_*^{-b} \int_\tau^t \lambda(\theta)^{-b} d\theta} \lambda(\tau) d\tau. \end{aligned}$$

(ii) Assume, in addition, that $M_0 \stackrel{\text{def}}{=} \sup_{[0, \infty)} \lambda(t) < \infty$. Then

$$\begin{aligned} \int_U |p_1(x, t) - p_2(x, t)|^2 dx &\leq e^{-C_1 t / (A_* M_0)^b} \int_U |p_1(x, 0) - p_2(x, 0)|^2 dx \\ &+ C_2 C_1^{-1} (A_* M_0)^{1+b} |\vec{a}^{(1)} - \vec{a}^{(2)}|. \end{aligned}$$

Consequently

$$\limsup_{t \rightarrow \infty} \int_U |p_1(x, t) - p_2(x, t)|^2 dx \leq C_2 C_1^{-1} (A_* M_0)^{1+b} |\vec{a}^{(1)} - \vec{a}^{(2)}|.$$

Forchheimer Equations in Porous Media - Part IV

- Boundary Flux condition
- Refined asymptotic estimates: limsup estimates for nonlinear differential inequalities, uniform Gronwall inequality, ...
- Continuous dependence for pressure gradients and velocity

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THANK YOU!