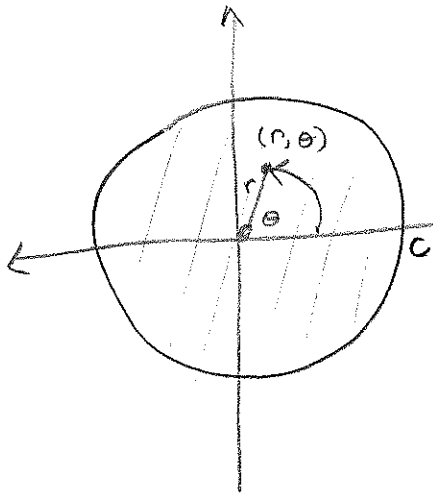


13.2 Wave Equation on a circular membrane

①



① We change the coordinates (x, y) to polar coordinates (r, θ) .

$$u(x, y, t) \Rightarrow u(r, \theta, t)$$

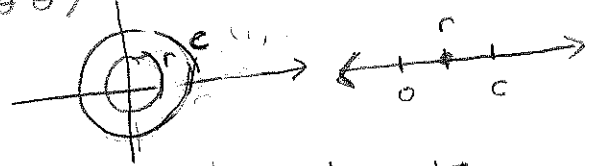
② We will do the case where the vibrations do not depend on the angular component θ :

$$u(r, \theta, t) \Rightarrow \boxed{u(r, t)}$$

Wave equation in polar coordinates:

(*) $u_{tt} = a^2 \left(u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta} \right)$ but $u_{\theta\theta} = 0$

(**) $\boxed{u_{tt} = a^2 \left(u_{rr} + \frac{1}{r} u_r \right)}$



We will discuss (*) with $u_{\theta\theta} \neq 0$ later.

Use separation of variables for (**),

$$\boxed{u(r, t) = R(r)T(t) = RT}$$

PDE: $u_{tt} = a^2 \left(u_{rr} + \frac{1}{r} u_r \right) \quad 0 < r < c, t > 0$

B.C. $u(c, t) = 0, t > 0$

I.C. $u(r, 0) = f(r), \left. \frac{\partial u}{\partial t} \right|_{t=0} = g(r), 0 < r < c$

$$u_{tt} = a^2 \left(u_{rr} + \frac{1}{r} u_r \right) \Rightarrow \frac{RT''}{a^2 RT} = \frac{a^2 \left(R''T + \frac{1}{r} R'T \right)}{a^2 RT} \quad (2)$$

$$\Rightarrow \frac{T''}{a^2 T} = \frac{R''}{R} + \frac{1}{r} \frac{R'}{R} = -\lambda$$

Two ODEs: $\underline{T'' + a^2 \lambda T = 0}$ and $\underline{R'' + \frac{1}{r} R' + \lambda R = 0}$

Rewrite the ODE in R to give an eigenvalue problem with the homogeneous B.C. $u(c, t) = 0$

$$(**) \quad \boxed{r^2 R'' + r R' + \lambda r^2 R = 0}$$

$$R(c) = 0$$

$\lim_{r \rightarrow 0} R(r)$ is bounded

The ODE $(**)$ is a Bessel's Equation of order zero. In general, a Bessel's Equation of order n is

$$\boxed{r^2 R'' + r R' + (\lambda r^2 - n^2) R = 0}$$

provided $\lambda > 0$.

For $\lambda = \alpha^2$, the general solution of $(**)$ is

$$R(r) = c_1 J_0(\alpha r) + c_2 Y_0(\alpha r)$$

$J_0(\alpha r)$ and $Y_0(\alpha r)$ are Bessel functions of order zero, first kind, and second kind, respectively. But $Y_0(\alpha r)$ is not bounded at $r=0$

$$\boxed{R(r) = J_0(\alpha r)}$$

Apply the BC $R(c) = 0$ to find the eigenvalues ⁽³⁾

$R(c) = J_0(\alpha c) = 0$ for $\alpha_1 < \alpha_2 < \alpha_3 < \dots < \alpha_n < \dots$

$\lambda_n = \alpha_n^2$ are the eigenvalues and

$R_n(r) = J_0(\alpha_n r)$ are the eigenfunctions,

From Sturm-Liouville Theory, these eigenfunctions are orthogonal on $[0, c]$ with respect to the weight function r

$$\int_0^c r J_0(\alpha_n r) J_0(\alpha_m r) dr = 0, n \neq m$$

Time equation: $T_n'' + \alpha_n^2 T = 0$

$$T_n(t) = a_n \cos(\alpha_n t) + b_n \sin(\alpha_n t)$$

Product solutions $u_n(r, t) = R_n(r) T_n(t)$

$$u(r, t) = \sum_{n=1}^{\infty} u_n(r, t)$$

$$= \sum_{n=1}^{\infty} \underbrace{[a_n \cos(\alpha_n t) + b_n \sin(\alpha_n t)]}_{T_n(t)} \underbrace{J_0(\alpha_n r)}_{R_n(r)}$$

We use properties of Bessel functions to compute $\alpha_n, n = 1, 2, 3, \dots$

Use of Computers

Solve PDE $u_{tt} = u_{rr} + \frac{1}{r}u_r$, $0 < r < 1$
 $t > 0$

B.C $u(1, t) = 0$, $t > 0$

I.C. $u(r, 0) = 0$, $0 < r < 1$

f.c. $\frac{\partial u}{\partial t} \Big|_{t=0} = \begin{cases} -1, & 0 \leq r < \frac{1}{4} \\ 0, & \frac{1}{4} < r < 1 \end{cases} = g(r)$

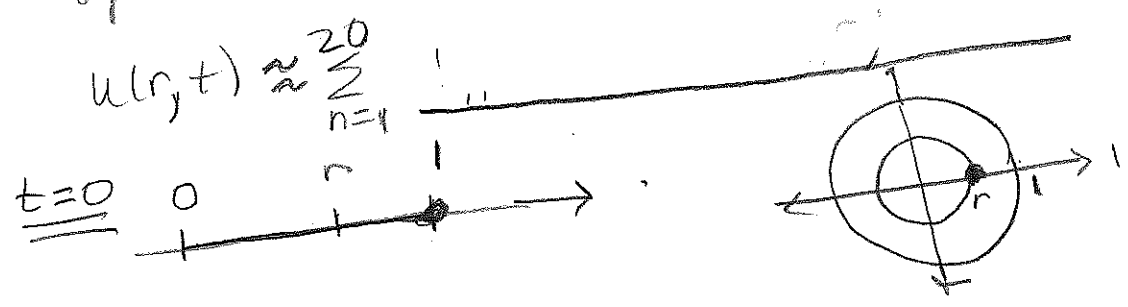
$$u(r, t) = \sum_{n=1}^{\infty} [a_n \cos(\alpha_n t) + b_n \sin(\alpha_n t)] J_0(\alpha_n r)$$

$u(r, 0) =$ _____ (coefficients)

$\frac{\partial u}{\partial t} \Big|_{t=0} = \sum_{n=1}^{\infty} b_n \alpha_n J_0(\alpha_n r) = g(r)$

$b_n \alpha_n = \frac{\int_0^1 r J_0(\alpha_n r) g(r) dr}{\int_0^1 r J_0^2(\alpha_n r) dr} = \frac{\int_0^{\frac{1}{4}} r J_0(\alpha_n r) dr}{\int_0^1 r J_0^2(\alpha_n r) dr}$

We use the computer to compute α_n and b_n up to $n=20$, and graph



Generalization of the Wave Equation

PDE $u_{tt} = a^2 \left(u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta} \right)$ $0 < r < c$
 $0 < \theta < 2\pi$

B.C. $u(c, \theta, t) = 0$
 $u(c, \theta, t) = u(c, \theta + 2\pi, t)$
 $t > 0$

I.C. $u(r, \theta, 0) = f(r, \theta)$
 $\frac{\partial u}{\partial t} \Big|_{t=0} = g(r, \theta)$

$u(r, \theta, t) = R(r) \Theta(\theta) T(t)$

$$\begin{cases} \Theta'' + \nu \Theta = 0 & \Theta(\theta) = \Theta(\theta + 2\pi) \\ r^2 R'' + r R' + (\lambda r^2 - \nu) R = 0, & R(c) = 0 \\ T'' + a^2 \lambda T = 0 \end{cases}$$

$\nu = n^2, \Theta_n = A_n \cos(n\theta) + B_n \sin(n\theta) \quad n=1, 2, \dots$
 $\lambda_{n,m} = \alpha_{n,m}^2, R_{n,m} = J_n(\alpha_{n,m} r), \quad n=1, 2, \dots, m=1, 2, \dots$

$T_{n,m} = C_{nm} \cos(\alpha_{n,m} t) + D_{nm} \sin(\alpha_{n,m} t)$
 $n=1, 2, \dots$
 $m=1, 2, \dots$

$u(r, \theta, t) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} T_{n,m}(t) R_{n,m}(r) \Theta_n(\theta)$

See problem # 17 p. 514

6

Question: If the initial position of the circular membrane is $u(r, \theta, 0) = 0$, write the simplified form of the solution to the wave equation.

$$u(r, \theta, t) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \underbrace{[C_{n,m} \cos(\alpha_{n,m} t) + D_{n,m} \sin(\alpha_{n,m} t)]}_{T_{n,m}} \underbrace{J_n(\alpha_{n,m} r)}_{R_{n,m}} \underbrace{[A_n \cos(n\theta) + B_n \sin(n\theta)]}_{\Theta_n}$$

$$u(r, \theta, 0) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty}$$

Coefficients

$$u(r, \theta, t) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty}$$

... Laplace Transform of $f(t)$ for $t \in [0, \infty)$.

$$* \mathcal{L}\{f(t)\} = \int_0^{\infty} e^{-st} f(t) dt = F(s)$$

$$\mathcal{L}\{1\} = \int_0^{\infty} e^{-st} 1 dt = \frac{1}{s}$$

$$\mathcal{L}\{e^{at}\} = \int_0^{\infty} e^{-st} e^{at} dt = \frac{1}{s-a}$$

$$\mathcal{L}\{\cos(at)\} = \frac{s}{s^2+a^2}$$

$$\mathcal{L}\{\sin(at)\} = \frac{a}{s^2+a^2}$$

We will use the identities for the derivative and apply Laplace transforms to solve PDEs with time derivative s .

$$\mathcal{L}\{f'(t)\} = s\mathcal{L}\{f(t)\} - f(0) = sF(s) - f(0)$$

$$\begin{aligned} \mathcal{L}\{f''(t)\} &= s^2\mathcal{L}\{f(t)\} - sf(0) - f'(0) \\ &= s^2F(s) - sf(0) - f'(0) \end{aligned}$$

Example $\mathcal{L}\{10 - 3\cos(2t)\} =$

$$\mathcal{L}^{-1}\left\{\frac{4}{s} + \frac{51}{s^2+4}\right\} =$$
