

Analysis of the Navier-Stokes systems and generalized Forchheimer flows

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Colloquium

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- ① The Navier-Stokes equations of viscous, incompressible fluids
 - ② Generalized Forchheimer flows in porous media
 - ③ Regularity theory for partial differential equations.
 - ④ Abstract dynamical systems.
- ▶ 29 articles published, 1 in press, 1 accepted, 1 submitted.
 - ▶ All are mathematical analysis, new results (except for one survey).
 - ▶ All are in peer-reviewed journals.

- 1 The Navier-Stokes systems
 - Foias-Saut asymptotic expansion
 - Exponentially decaying forces
 - Power-decaying forces
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1. The Navier-Stokes systems

- Foias-Saut asymptotic expansion
- Exponentially decaying forces
- Power-decaying forces

The Navier-Stokes equations

- The Navier-Stokes equations (NSE) in \mathbb{R}^3 :

$$\begin{cases} \frac{\partial u}{\partial t} + (u \cdot \nabla)u - \nu \Delta u + \nabla p = f(x, t), \\ \operatorname{div} u = 0, \\ u(x, 0) = u^0(x), \end{cases}$$

with viscosity $\nu > 0$, velocity field $u(x, t) \in \mathbb{R}^3$, pressure $p(x, t) \in \mathbb{R}$, body force $f(x, t) \in \mathbb{R}^3$, initial velocity $u^0(x)$.

- Let $L > 0$ and $\Omega = (0, L)^3$. The L-periodic solutions:

$$u(x + Le_j) = u(x) \text{ for all } x \in \mathbb{R}^3, j = 1, 2, 3,$$

where $\{e_1, e_2, e_3\}$ is the canonical basis in \mathbb{R}^3 .

Zero average condition

$$\int_{\Omega} u(x) dx = 0,$$

Throughout $L = 2\pi$ and $\nu = 1$.

Functional setting

Let \mathcal{V} be the set of \mathbb{R}^3 -valued 2π -periodic trigonometric polynomials which are divergence-free and satisfy the zero average condition.

$$H = \text{closure of } \mathcal{V} \text{ in } L^2(\Omega)^3 = H^0(\Omega)^3,$$

$$V = \text{closure of } \mathcal{V} \text{ in } H^1(\Omega)^3, \quad \mathcal{D}(A) = \text{closure of } \mathcal{V} \text{ in } H^2(\Omega)^3.$$

Norm on H : $|u| = \|u\|_{L^2(\Omega)}$. Norm on V : $\|u\| = |\nabla u|$.

The Stokes operator:

$$Au = -\Delta u \text{ for all } u \in \mathcal{D}(A).$$

The bilinear mapping:

$$B(u, v) = \mathbb{P}_L(u \cdot \nabla v) \text{ for all } u, v \in \mathcal{D}(A).$$

\mathbb{P}_L is the Leray projection from $L^2(\Omega)$ onto H .

WLOG, assume $f(t) = \mathbb{P}_L f(t)$. The functional form of the NSE:

$$\frac{du(t)}{dt} + Au(t) + B(u(t), u(t)) = f(t), \quad t > 0,$$

$$u(0) = u^0.$$

Case $f = 0$. Foias-Saut asymptotic expansion

Foias-Saut (1987) for a solution $u(t)$:

$$u(t) \sim \sum_{n=1}^{\infty} q_n(t) e^{-jt},$$

where $q_j(t)$ is a \mathcal{V} -valued polynomial in t . This means that for any $N \in \mathbb{N}$, $m \in \mathbb{N}$, the remainder $v_N(t) = u(t) - \sum_{j=1}^N q_j(t) e^{-jt}$ satisfies

$$\|v_N(t)\|_{H^m(\Omega)} = O(e^{-(N+\varepsilon)t})$$

as $t \rightarrow \infty$, for some $\varepsilon = \varepsilon_{N,m} > 0$.

Theorem (H.-Martinez 2017)

The Foias-Saut expansion holds in all Gevrey spaces:

$$\|e^{\sigma A^{1/2}} v_N(t)\|_{H^m(\Omega)} = O(e^{-(N+\varepsilon)t}),$$

for any $\sigma > 0$, $\varepsilon \in (0, 1)$.

Gevrey classes

- Spectrum of A is $\{|k|^2 : k \in \mathbb{Z}^3, k \neq 0\}$.
- For $\alpha \geq 0, \sigma \geq 0$, define

$$A^\alpha e^{\sigma A^{1/2}} u = \sum_{\mathbf{k} \neq 0} |\mathbf{k}|^{2\alpha} \hat{u}(\mathbf{k}) e^{\sigma |\mathbf{k}|} e^{i\mathbf{k} \cdot \mathbf{x}}, \text{ for } u = \sum_{\mathbf{k} \neq 0} \hat{u}(\mathbf{k}) e^{i\mathbf{k} \cdot \mathbf{x}} \in H.$$

The domain of $A^\alpha e^{\sigma A^{1/2}}$ is

$$G_{\alpha, \sigma} = \mathcal{D}(A^\alpha e^{\sigma A^{1/2}}) = \{u \in H : |u|_{\alpha, \sigma} \stackrel{\text{def}}{=} |A^\alpha e^{\sigma A^{1/2}} u| < \infty\}.$$

- Compare the Sobolev and Gevrey norms:

$$|A^\alpha u| = |(A^\alpha e^{-\sigma A^{1/2}}) e^{\sigma A^{1/2}} u| \leq \left(\frac{2\alpha}{e\sigma}\right)^{2\alpha} |e^{\sigma A^{1/2}} u|.$$

Notation.

- Denote for $\sigma \in \mathbb{R}$ the space $E^{\infty, \sigma} = \bigcap_{\alpha \geq 0} G_{\alpha, \sigma} = \bigcap_{m \in \mathbb{N}} G_{m, \sigma}$.
- Denote by $\mathcal{P}^{\alpha, \sigma}$ the space of $G_{\alpha, \sigma}$ -valued polynomials in case $\alpha \in \mathbb{R}$, and the space of $E^{\infty, \sigma}$ -valued polynomials in case $\alpha = \infty$.

Definition

Let X be a real vector space.

(a) An X -valued polynomial is a function $t \in \mathbb{R} \mapsto \sum_{n=1}^d a_n t^n$, for some $d \geq 0$, and a_n 's belonging to X .

(b) In case $\|\cdot\|$ is a norm on X , a function $g(t)$ from $(0, \infty)$ to X is said to have the asymptotic expansion

$$g(t) \sim \sum_{n=1}^{\infty} g_n(t) e^{-nt} \text{ in } X,$$

where $g_n(t)$'s are X -valued polynomials, if for all $N \geq 1$, there exists $\varepsilon_N > 0$ such that

$$\left\| g(t) - \sum_{n=1}^N g_n(t) e^{-nt} \right\| = \mathcal{O}(e^{-(N+\varepsilon_N)t}) \text{ as } t \rightarrow \infty.$$

• We will say that an asymptotic expansion holds in $E^{\infty, \sigma}$ if it holds in $G_{\alpha, \sigma}$ for all $\alpha \geq 0$.

Exponentially decaying forces

(A1) The function $f(t)$ is continuous from $[0, \infty)$ to H .

(A2) There are a number $\sigma_0 \geq 0$, E^{∞, σ_0} -valued polynomials $f_n(t)$ such that

$$f(t) \sim \sum_{n=1}^{\infty} f_n(t)e^{-nt} \text{ in } E^{\infty, \sigma_0}.$$

Theorem (H.-Martinez 2018)

Let $u(t)$ be a Leray-Hopf weak solution. Then $u(t)$ has the asymptotic expansion

$$u(t) \sim \sum_{n=1}^{\infty} q_n(t)e^{-nt} \text{ in } E^{\infty, \sigma_0}.$$

Moreover, $u_n(t) \stackrel{\text{def}}{=} q_n(t)e^{-nt}$ and $F_n(t) \stackrel{\text{def}}{=} f_n(t)e^{-nt}$ satisfy

$$\frac{d}{dt}u_n(t) + Au_n(t) + \sum_{\substack{k, m \geq 1 \\ k+m=n}} B(u_k(t), u_m(t)) = F_n(t), \quad t \in \mathbb{R}, \text{ in } E^{\infty, \sigma_0}.$$

Finite asymptotic approximation

Theorem (H.-Martinez 2018)

Suppose there exist $N_* \geq 1$, $\sigma_0 \geq 0$, $\mu_* \geq \alpha_* \geq N_*/2$, and $\delta_n \in (0, 1)$ and $f_n \in \mathcal{P}^{\mu_n, \sigma_0}$, for $1 \leq n \leq N_*$, such that

$$\left| f(t) - \sum_{n=1}^N f_n(t) e^{-nt} \right|_{\alpha_N, \sigma_0} = \mathcal{O}(e^{-(N+\delta_N)t}) \quad \text{as } t \rightarrow \infty,$$

for $1 \leq N \leq N_*$, where $\mu_n = \mu_* - (n-1)/2$, $\alpha_n = \alpha_* - (n-1)/2$.

Let $u(t)$ be a Leray-Hopf weak . Then there exist polynomials $q_n \in \mathcal{P}^{\mu_n+1, \sigma_0}$, for $1 \leq n \leq N_*$, such that one has for $1 \leq N \leq N_*$ that

$$\left| u(t) - \sum_{n=1}^N q_n(t) e^{-nt} \right|_{\alpha_N, \sigma_0} = \mathcal{O}(e^{-(N+\varepsilon)t}) \quad \text{as } t \rightarrow \infty, \quad \forall \varepsilon \in (0, \delta_N^*),$$

where $\delta_N^* = \min\{\delta_1, \delta_2, \dots, \delta_N\}$.

Power-decaying forces

Power asymptotic expansion in $(X, \|\cdot\|)$: $g(t) \stackrel{\text{pow.}}{\sim} \sum_{n=1}^{\infty} g_n t^{-n}$ means

$$\|g(t) - \sum_{n=1}^N g_n t^{-n}\| = \mathcal{O}(t^{-(N+\varepsilon)}), \quad \text{for some } \varepsilon > 0, \quad t \rightarrow \infty.$$

Theorem (Cao-H. 2017)

Assume that $f(t) \stackrel{\text{pow.}}{\sim} \sum_{n=1}^{\infty} \phi_n t^{-n}$ in G_{α, σ_0} , for some $\sigma_0 \geq 0$ and $\alpha \geq 1/2$, sequence $\{\phi_n\}_{n=1}^{\infty}$ in G_{α, σ_0} . Then any Leray-Hopf weak solution $u(t)$ has the asymptotic expansion

$$u(t) \stackrel{\text{pow.}}{\sim} \sum_{n=1}^{\infty} \xi_n t^{-n} \quad \text{in } G_{\alpha, \sigma_0},$$

where $\xi_1 = A^{-1}\phi_1$,

$\xi_n = (n-1)A^{-1}\xi_{n-1} - \sum_{k,m \geq 1, k+m=n} A^{-1}B(\xi_k, \xi_m) + A^{-1}\phi_n$ for $n \geq 2$.

2. Expansions in a general system of decaying functions

- Continuum systems
- Asymptotic expansions for NSE
- Applications

Expansions in a general system of decaying functions

Definition (Very/Too general)

Let $(\psi_n)_{n=1}^{\infty}$ be a sequence of non-negative functions defined on $[T_*, \infty)$ for some $T_* \in \mathbb{R}$ that satisfies the following two conditions:

- 1 For each $n \in \mathbb{N}$, $\lim_{t \rightarrow \infty} \psi_n(t) = 0$.
- 2 For $n > m$, $\psi_n(t) = o(\psi_m(t))$.

Let $(X, \|\cdot\|)$ be a normed space, and g be a function from $[T_*, \infty)$ to X .

$$g(t) \sim \sum_{n=1}^{\infty} \xi_n \psi_n(t) \text{ in } X,$$

where $\xi_n \in X$ for all $n \in \mathbb{N}$, if, for any $N \in \mathbb{N}$,

$$\|g(t) - \sum_{n=1}^N \xi_n \psi_n(t)\| = o(\psi_N(t)).$$

Definition

Let $\Psi = (\psi_\lambda)_{\lambda>0}$ be a system of functions that satisfies the following two conditions.

- (a) There exists $T_* \geq 0$ such that, for each $\lambda > 0$, ψ_λ is a positive function defined on $[T_*, \infty)$, and

$$\lim_{t \rightarrow \infty} \psi_\lambda(t) = 0.$$

- (b) For any $\lambda > \mu$, there exists $\eta > 0$ such that

$$\psi_\lambda(t) = \mathcal{O}(\psi_\mu(t)\psi_\eta(t)).$$

Definition

Let $(X, \|\cdot\|)$ be a real normed space, and $g : (0, \infty) \rightarrow X$.

$$g(t) \overset{\Psi}{\sim} \sum_{n=1}^{\infty} \xi_n \psi_{\lambda_n}(t) \text{ in } X,$$

where $\xi_n \in X$ for all $n \in \mathbb{N}$, and $(\lambda_n)_{n=1}^{\infty}$ is a strictly increasing, divergent sequence of positive numbers, if it holds, for any $N \geq 1$, that there exists $\varepsilon > 0$ such that

$$\left\| g(t) - \sum_{n=1}^N \xi_n \psi_{\lambda_n}(t) \right\| = \mathcal{O}(\psi_{\lambda_N}(t) \psi_{\varepsilon}(t)).$$

Condition

The system $\Psi = (\psi_\lambda)_{\lambda>0}$ satisfies (a) and (b) and the following.

- 1 For any $\lambda, \mu > 0$, there exist $\gamma > \max\{\lambda, \mu\}$ and a nonzero constant $d_{\lambda,\mu}$ such that

$$\psi_\lambda \psi_\mu = d_{\lambda,\mu} \psi_\gamma.$$

Notation. $\gamma = \lambda \wedge \mu$.

- 2 For each $\lambda > 0$, the function ψ_λ is continuous and differentiable on $[T_*, \infty)$, and its derivative ψ'_λ has an expansion

$$\psi'_\lambda(t) \stackrel{\Psi}{\sim} \sum_{k=1}^{N_\lambda} c_{\lambda,k} \psi_{\lambda^\vee(k)}(t) \text{ in } \mathbb{R},$$

where $N_\lambda \in \mathbb{N} \cup \{0, \infty\}$, all $c_{\lambda,k}$ are constants, all $\lambda^\vee(k) > \lambda$, and, for each $\lambda > 0$, $\lambda^\vee(k)$'s are strictly increasing in k .

Condition

The system $\Psi = (\psi_\lambda)_{\lambda>0}$ satisfies (a), (b) and the following.

- 1 For each $\lambda > 0$, the function ψ_λ is decreasing (in t).
- 2 If $\lambda, \alpha > 0$ then

$$e^{-\alpha t} = o(\psi_\lambda(t)).$$

- 3 For any number $a \in (0, 1)$,

$$\psi_\lambda(at) = \mathcal{O}(\psi_\lambda(t)).$$

Consequently, for any $T \in \mathbb{R}$,

$$\psi_\lambda(t) = \mathcal{O}(\psi_\lambda(t + T)).$$

Assumption

The function f belongs to $L_{\text{loc}}^\infty([0, \infty), H)$.

Asymptotic expansions for NSE

Assumption

Suppose there exist real numbers $\sigma \geq 0$, $\alpha \geq 1/2$, a strictly increasing, divergent sequence of positive numbers $(\gamma_n)_{n=1}^{\infty}$ and a sequence $(\tilde{\phi}_n)_{n=1}^{\infty}$ in $G_{\alpha,\sigma}$ such that

$$f(t) \stackrel{\Psi}{\sim} \sum_{n=1}^{\infty} \tilde{\phi}_n \psi_{\gamma_n}(t) \text{ in } G_{\alpha,\sigma}.$$

Assumption

There exists a set S_* that contains $\{\gamma_n : n \in \mathbb{N}\}$, preserves the operations \vee and \wedge , and can be ordered so that

$S_* = \{\lambda_n : n \in \mathbb{N}\}$, where λ_n 's are strictly increasing to infinity.

Rewrite

$$f(t) \stackrel{\Psi}{\sim} \sum_{n=1}^{\infty} \phi_n \psi_{\lambda_n}(t) \text{ in } G_{\alpha,\sigma} \text{ as } t \rightarrow \infty,$$

Theorem (Cao-H. 2018)

Any Leray-Hopf weak solution $u(t)$ has the asymptotic expansion

$$u(t) \overset{\Psi}{\sim} \sum_{n=1}^{\infty} \xi_n \psi_{\lambda_n}(t) \quad \text{in } G_{\alpha+1-\rho, \sigma} \text{ for all } \rho \in (0, 1),$$

where ξ_n 's are defined recursively by

$$\xi_1 = A^{-1} \phi_1,$$

$$\xi_n = A^{-1} \left(\phi_n - \chi_n - \sum_{\substack{1 \leq k, m \leq n-1, \\ \lambda_k \wedge \lambda_m = \lambda_n}} d_{\lambda_k, \lambda_m} B(\xi_k, \xi_m) \right) \quad \text{for } n \geq 2,$$

where

$$\chi_n = \begin{cases} \sum_{\substack{(p, k) \in [1, n-1] \times \mathbb{N}: \\ \lambda_p^\vee(k) = \lambda_n}} c_{\lambda_p, k} \xi_p, & \text{if } \exists p \in [1, n-1], k \in \mathbb{N} : \lambda_p^\vee(k) = \lambda_n, \\ 0, & \text{otherwise.} \end{cases}$$

Theorem (Cao-H. 2018)

Given $\alpha, \sigma \geq 0$, let $\xi \in G_{\alpha, \sigma}$, and f be a function from $(0, \infty)$ to $G_{\alpha, \sigma}$ that satisfies

$$|f(t)|_{\alpha, \sigma} \leq MF(t) \quad \text{a.e. in } (0, \infty),$$

where F is a continuous, decreasing function from $[0, \infty)$ to $[0, \infty)$. Let $w_0 \in G_{\alpha, \sigma}$. Suppose $w \in C([0, \infty), H_w) \cap L^1_{\text{loc}}([0, \infty), V)$, with $w' \in L^1_{\text{loc}}([0, \infty), V')$, is a weak solution of

$$w' = -Aw + \xi + f \text{ in } V' \text{ on } (0, \infty), \quad w(0) = w_0,$$

Then the following statements hold true.

- 1 $w(t) \in G_{\alpha+1-\varepsilon, \sigma}$ for all $\varepsilon \in (0, 1)$ and $t > 0$.

Theorem (continued)

- ② For any numbers $a, a_0 \in (0, 1)$ with $a + a_0 < 1$ and any $\varepsilon \in (0, 1)$, there exists a positive constant C depending on $a_0, a, \varepsilon, M, F(0), |\xi|_{\alpha, \sigma}$ and $|w_0|_{\alpha, \sigma}$ such that

$$|w(t) - A^{-1}\xi|_{\alpha+1-\varepsilon, \sigma} \leq C(e^{-a_0 t} + F(at)) \quad \forall t \geq 1.$$

- ③ Assume, in addition, that

- There exist $k_0 > 0$ and $D_1 > 0$ such that

$$e^{-k_0 t} \leq D_1 F(t) \quad \forall t \geq 0, \text{ and} \quad (\text{F1})$$

- For any $a \in (0, 1)$, there exists $D_2 = D_{2,a} > 0$ such that

$$F(at) \leq D_2 F(t) \quad \forall t \geq 0. \quad (\text{F2})$$

Then there exists $C > 0$ such that

$$|w(t) - A^{-1}\xi|_{\alpha+1-\varepsilon, \sigma} \leq CF(t) \quad \forall t \geq 1.$$

Theorem (Cao-H. 2018)

Let F be a continuous, decreasing, non-negative function on $[0, \infty)$. Given $\alpha \geq 1/2$ and numbers $\theta_0, \theta \in (0, 1)$ such that $\theta_0 + \theta < 1$. Then there exist positive numbers $c_k = c_k(\alpha, \theta_0, \theta, F)$, for $k = 0, 1, 2, 3$, such that the following holds true. If

$$\begin{aligned} |A^\alpha u^0| &\leq c_0, \\ |f(t)|_{\alpha-1/2, \sigma} &\leq c_1 F(t) \quad \text{a.e. in } (0, \infty) \text{ for some } \sigma \geq 0, \end{aligned}$$

then there exists a unique regular solution $u(t)$, which also belongs to $C([0, \infty), \mathcal{D}(A^\alpha))$ and satisfies, for all $t \geq 8\sigma(1-\theta)/(1-\theta-\theta_0)$,

$$\begin{aligned} |u(t)|_{\alpha, \sigma} &\leq c_2(e^{-2\theta_0 t} + F^2(\theta t))^{1/2}, \\ \int_t^{t+1} |u(\tau)|_{\alpha+1/2, \sigma}^2 d\tau &\leq c_3^2(e^{-2\theta_0 t} + F^2(\theta t)). \end{aligned}$$

Parts of proof (3). Estimates for Leray-Hopf weak solutions

Theorem (Cao-H. 2018)

Let F be a continuous, decreasing, non-negative function such that

$$\lim_{t \rightarrow \infty} F(t) = 0,$$

$$|f(t)|_{\alpha, \sigma} = \mathcal{O}(F(t)), \text{ for some } \sigma \geq 0, \alpha \geq 1/2.$$

Let $u(t)$ be a Leray-Hopf weak solution. Then there exists $\hat{T} > 0$ such that $u(t)$ is a regular solution on $[\hat{T}, \infty)$, and for any $\varepsilon, \lambda \in (0, 1)$, and $a_0, a, \theta_0, \theta \in (0, 1)$ with $a_0 + a < 1$, $\theta_0 + \theta < 1$,

$$|u(\hat{T} + t)|_{\alpha+1-\varepsilon, \sigma} \leq C(e^{-a_0 t} + e^{-2\theta_0 a t} + F^{2\lambda}(\theta a t) + F(at)) \quad \forall t \geq 0.$$

If, in addition, F satisfies (F1) and (F2), then

$$|u(\hat{T} + t)|_{\alpha+1-\varepsilon, \sigma} \leq CF(t) \quad \forall t \geq 0.$$

Application: iterated logarithmic decaying functions

For $k, m \in \mathbb{N}$, let

$$L_k(t) = \underbrace{\ln(\ln(\cdots \ln(t)))}_{k\text{-times}} \quad \text{and} \quad \mathcal{L}_m(t) = (L_1(t), L_2(t), \cdots, L_m(t)).$$

- Let $Q_0 : \mathbb{R}^m \rightarrow \mathbb{R}$ be a polynomial in m variables with positive degree and positive leading coefficient:

$$Q_0(z) = \sum_{\alpha} c_{\alpha} z^{\alpha} \quad \text{for } z \in \mathbb{R}^m.$$

We use the lexicographic order for the multi-indices.

- Let Q_1 be a polynomial in one variable of positive degree with positive leading coefficient.

Given a number $\beta > 0$, we define

$$\omega(t) = (Q_0 \circ \mathcal{L}_m \circ Q_1)(t^{\beta}) \quad \text{with } t \in \mathbb{R}.$$

Let $\psi_\lambda(t) = \omega(t)^{-\lambda}$ and $\Psi = (\psi_\lambda(t))_{\lambda>0}$. Note $\psi'_\lambda \stackrel{\Psi}{\sim} 0$.

Theorem (Cao-H. 2018)

Assume

$$f(t) \stackrel{\Psi}{\sim} \sum_{n=1}^{\infty} \phi_n \omega(t)^{-\lambda_n} \quad \text{in } G_{\alpha,\sigma},$$

for some $\sigma \geq 0$, $\alpha \geq 1/2$. Then any Leray-Hopf weak solution $u(t)$ of the NSE has the asymptotic expansion

$$u(t) \stackrel{\Psi}{\sim} \sum_{n=1}^{\infty} \xi_n \omega(t)^{-\lambda_n} \quad \text{in } G_{\alpha+1-\rho,\sigma} \quad \text{for all } \rho \in (0, 1),$$

where

$$\xi_1 = A^{-1}\phi_1, \quad \xi_n = A^{-1}\left(\phi_n - \sum_{\substack{1 \leq k, m \leq n-1, \\ \lambda_k + \lambda_m = \lambda_n}} B(\xi_k, \xi_m)\right) \quad \text{for } n \geq 2.$$

Corollary (Cao-H. 2018)

Given $m \in \mathbb{N}$, define $\Psi = (L_m(t)^{-\lambda})_{\lambda > 0}$. Suppose $(\lambda_n)_{n=1}^{\infty}$ is a strictly increasing, divergent sequence of positive numbers such that the set $\{\lambda_n : n \in \mathbb{N}\}$ preserves the addition. If

$$f(t) \underset{\Psi}{\sim} \sum_{n=1}^{\infty} \phi_n L_m(t)^{-\lambda_n} \quad \text{in } G_{\alpha, \sigma},$$

then any Leray-Hopf weak solution $u(t)$ of the NSE admits the asymptotic expansion

$$u(t) \underset{\Psi}{\sim} \sum_{n=1}^{\infty} \xi_n L_m(t)^{-\lambda_n} \quad \text{in } G_{\alpha+1-\rho, \sigma} \text{ for all } \rho \in (0, 1).$$

Expansions with trigonometric functions

Example. If

$$f(t) \underset{\Psi}{\sim} \sum_{n=1}^{\infty} \phi_n [\sin(L_m^{-1}(t))]^{\lambda_n} \quad \text{in } G_{\alpha,\sigma},$$

then

$$u(t) \underset{\Psi}{\sim} \sum_{n=1}^{\infty} \xi_n [\sin(L_m^{-1}(t))]^{\lambda_n} \quad \text{in } G_{\alpha+1-\rho,\sigma} \text{ for all } \rho \in (0, 1),$$

Example. If

$$f(t) \underset{\Psi}{\sim} \sum_{n=1}^{\infty} \phi_n [\tan(L_m^{-1}(t))]^{\lambda_n} \quad \text{in } G_{\alpha,\sigma},$$

then

$$u(t) \underset{\Psi}{\sim} \sum_{n=1}^{\infty} \xi_n [\tan(L_m^{-1}(t))]^{\lambda_n} \quad \text{in } G_{\alpha+1-\rho,\sigma} \text{ for all } \rho \in (0, 1),$$

Infinite expansions for the derivatives

Consider $\Psi = (\psi_\lambda)_{\lambda>0}$ with $\psi_\lambda = (\sqrt{t} + 1)^{-\lambda}$. Then

$$\begin{aligned}\psi'_\lambda(t) &= -\lambda(\sqrt{t} + 1)^{-\lambda-1} \frac{1}{2} \frac{1}{\sqrt{t}} = -\frac{\lambda}{2}(\sqrt{t} + 1)^{-\lambda-1} \frac{1}{\sqrt{t} + 1} \cdot \frac{1}{1 - \frac{1}{\sqrt{t}+1}} \\ &= \sum_{k=1}^{\infty} -\frac{\lambda}{2}(\sqrt{t} + 1)^{-\lambda-k-1}.\end{aligned}$$

Proposition (Cao-H. 2018)

Assume $f(t) \stackrel{\Psi}{\sim} \sum_{n=1}^{\infty} \phi_n(\sqrt{t} + 1)^{-\lambda_n}$ in $G_{\alpha,\sigma}$. Then

$$u(t) \stackrel{\Psi}{\sim} \sum_{n=1}^{\infty} \xi_n(\sqrt{t} + 1)^{-\lambda_n} \quad \text{in } G_{\alpha+1-\rho,\sigma} \text{ for all } \rho \in (0, 1),$$

where $\xi_1 = A^{-1}\phi_1$, $\xi_n = A^{-1}(\phi_n + \frac{1}{2} \sum_{p \in \mathcal{Z}_n} \lambda_p \xi_p - \sum_{\substack{1 \leq k, m \leq n-1, \\ \lambda_k + \lambda_m = \lambda_n}} B(\xi_k, \xi_m))$
for $n \geq 2$, with $\mathcal{Z}_n = \{p \in \mathbb{N} \cap [1, n-1] : \exists k \in \mathbb{N}, \lambda_p + 1 + k = \lambda_n\}$.

3. Generalized Forchheimer gas flows in porous media

- Generalized Forchheimer flows
- Gas flows: mathematical model
- Estimates of the Lebesgue norms
- Maximum estimates
- Gradient estimates

Darcy's and Forchheimer's flows

Fluid flows in porous media with velocity \mathbf{v} and pressure p :

- Darcy's Law:

$$\alpha \mathbf{v} = -\nabla p,$$

- Forchheimer's "two term" law

$$\alpha \mathbf{v} + \beta |\mathbf{v}| \mathbf{v} = -\nabla p,$$

- Forchheimer's "three term" law

$$\mathcal{A} \mathbf{v} + \mathcal{B} |\mathbf{v}| \mathbf{v} + \mathcal{C} |\mathbf{v}|^2 \mathbf{v} = -\nabla p.$$

- Forchheimer's "power" law

$$a \mathbf{v} + c^n |\mathbf{v}|^{n-1} \mathbf{v} = -\nabla p,$$

Here $\alpha, \beta, a, c, n, \mathcal{A}, \mathcal{B}$, and \mathcal{C} are empirical positive constants.

Generalized Forchheimer equations

[Aulisa-Bloshanskaya-H.-Ibragimov 2009]: $g(|v|)v = -\nabla p$.

Let $G(s) = sg(s)$. Then $G(|v|) = |\nabla p| \Rightarrow |v| = G^{-1}(|\nabla p|)$. Hence

$$v = -\frac{\nabla p}{g(G^{-1}(|\nabla p|))} \Rightarrow v = -K(|\nabla p|)\nabla p,$$

$$K(\xi) = K_g(\xi) = \frac{1}{g(s)} = \frac{1}{g(G^{-1}(\xi))}, \quad sg(s) = \xi.$$

Class $FP(N, \vec{\alpha})$. Let $N > 0$, $0 = \alpha_0 < \alpha_1 < \alpha_2 < \dots < \alpha_N$,

$$FP(N, \vec{\alpha}) = \left\{ g(s) = a_0 s^{\alpha_0} + a_1 s^{\alpha_1} + a_2 s^{\alpha_2} + \dots + a_N s^{\alpha_N} \right\}.$$

Lemma (Degeneracy)

$$\frac{C_1(\vec{a})}{(1 + \xi)^a} \leq K(\xi, \vec{a}) \leq \frac{C_2(\vec{a})}{(1 + \xi)^a},$$

$$C_3(\vec{a})(\xi^{2-a} - 1) \leq K(\xi, \vec{a})\xi^2 \leq C_2(\vec{a})\xi^{2-a}.$$

Works on Forchheimer flows

- Darcy-Dupuit: 1865
- Forchheimer: 1901
- Other nonlinear models: 1940s–1960s
- Incompressible fluids: Payne, Straughan and collaborators since 1990's, Celebi-Kalantarov-Ugurlu since 2005 (Brinkman-Forchheimer)
- Derivation of non-Darcy, non-Forchheimer flows: Marusic-Paloka and Mikelic 2009 (homogenization for Navier–Stokes equations), Balhoff et. al. 2009 (computational)

Works on generalized Forchheimer flows

- 1990's Numerical study
- L^2 -theory: Aulisa-Bloshanskaya-H.-Ibragimov (2009), H.-Ibragimov: Dirichlet B.C. (2011), H.-Ibragimov Flux B.C. (2012), Aulisa-Bloshanskaya-Ibragimov total flux, productivity index (2011, 2012).
- L^α -theory: H.-Ibragimov-Kieu-Sobol (2015)
- $L^\infty, W^{1,p}$ -theory: H.-Kieu-Phan (2014).
- $W^{1,\infty}$ -theory: interior H.-Kieu (2017), global H.-Kieu (2018-in press).
- Heterogeneous porous media: Celik-H.(2016, 2017).
- **Isentropic gases:** Celik-H.-Kieu (2018a, 2018b).
- **Mixed pre-Darcy, Darcy, Forchheimer flows:** Celik-H.-Ibragimov-Kieu (2017)
- Two-phase flows: H.-Ibragimov-Kieu (2013, 2014)
- Numericals: Kieu (2016, 2017, 2018) Ibragimov-Kieu (2016)

Gas flows: mathematical model [Celik-H.-Kieu 2018]

Based on dimensional analysis by Muskat and Ward, we consider

$$\sum_{i=0}^N a_i \rho^{\alpha_i} |\mathbf{v}|^{\alpha_i} \mathbf{v} = -\nabla p + \rho \vec{\mathbf{g}}.$$

Multiplying both sides by ρ , we obtain

$$g(|\rho \mathbf{v}|) \rho \mathbf{v} = -\rho \nabla p + \rho^2 \vec{\mathbf{g}}.$$

Solving for $\rho \mathbf{v}$ gives

$$\rho \mathbf{v} = -K(|\rho \nabla p - \rho^2 \vec{\mathbf{g}}|)(\rho \nabla p - \rho^2 \vec{\mathbf{g}}),$$

where the function $K : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ is defined for $\xi \geq 0$ by

$$K(\xi) = \frac{1}{g(s(\xi))},$$

with $s(\xi) = s$ being the unique non-negative solution of $sg(s) = \xi$.

Doubly nonlinear parabolic equations

Conservation of mass:

$$\phi \rho_t + \operatorname{div}(\rho \mathbf{v}) = F,$$

where the porosity $\phi \in (0, 1)$, F is the *source term*. Then

$$\phi \rho_t = \operatorname{div}(K(|\rho \nabla p - \rho^2 \vec{\mathbf{g}}|)(\rho \nabla p - \rho^2 \vec{\mathbf{g}})) + F.$$

Isentropic gas flows. In this case

$$p = \bar{c} \rho^\gamma \quad \text{for some constants } \bar{c}, \gamma > 0.$$

Here, γ is the specific heat ratio. Note that $\rho \nabla p = \nabla(\bar{c} \gamma \rho^{\gamma+1}/(\gamma+1))$.

Let (pseudo-pressure) $u = \frac{\bar{c} \gamma \rho^{\gamma+1}}{\gamma+1} = \frac{\gamma p^{\frac{\gamma+1}{\gamma}}}{\bar{c}^{\frac{1}{\gamma}}(\gamma+1)}$, we have

$$\phi c^{1/2} (u^\lambda)_t = \nabla \cdot (K(|\nabla u - cu^\ell \vec{\mathbf{g}}|)(\nabla u - cu^\ell \vec{\mathbf{g}})) + F,$$

where $\lambda = \frac{1}{\gamma+1} \in (0, 1)$, $\ell = 2\lambda$ and $c = \left(\frac{\gamma+1}{\bar{c}\gamma}\right)^\ell$.

Ideal gases. The equation of state is

$$p = \bar{c}\rho \quad \text{for some constant } \bar{c} > 0.$$

Same equation with $\gamma = 1$, $\lambda = 1/2$, the pseudo-pressure $u \sim p^2$.

Slightly compressible fluids. The equation of state is

$$\frac{1}{\rho} \frac{d\rho}{dp} = \frac{1}{\kappa} = \text{const.} > 0.$$

Then $\rho \nabla p = \kappa \nabla \rho$. Same equation with

$$\lambda = 1, \quad u = \kappa \rho, \quad \ell = 2 \quad \text{and} \quad c = 1/\kappa^2.$$

Boundary condition. Volumetric flux condition

$$\nu \cdot \vec{\nu} = \psi \quad \text{on } \partial U.$$

This gives $\rho \nu \cdot \vec{\nu} = \psi \rho$, hence,

$$-K(|\nabla u - cu^\ell \vec{\mathbf{g}}|)(\nabla u - cu^\ell \vec{\mathbf{g}}) \cdot \vec{\nu} = c^{1/2} \psi u^\lambda.$$

General formulation and the initial boundary value problem

$$\begin{cases} \frac{\partial(u^\lambda)}{\partial t} = \nabla \cdot (K(|\nabla u + Z(u)|)(\nabla u + Z(u))) + f(x, t, u) & \text{on } U \times (0, \infty), \\ u(x, 0) = u_0(x) & \text{on } U, \\ K(|\nabla u + Z(u)|)(\nabla u + Z(u)) \cdot \vec{\nu} = B(x, t, u) & \text{on } \Gamma \times (0, \infty), \end{cases}$$

where $u(x, t) \geq 0$

Assumption (A1). Assume $Z(u) : [0, \infty) \rightarrow \mathbb{R}^n$,

$B(x, t, u) : \Gamma \times [0, \infty) \times [0, \infty) \rightarrow \mathbb{R}$, and

$f(x, t, u) : U \times [0, \infty) \times [0, \infty) \rightarrow \mathbb{R}$ satisfy

$$|Z(u)| \leq d_0 u^{\ell_Z},$$

$$B(x, t, u) \leq \varphi_1(x, t) + \varphi_2(x, t) u^{\ell_B},$$

$$f(x, t, u) \leq f_1(x, t) + f_2(x, t) u^{\ell_f}$$

with constants $d_0, \ell_Z > 0$, $\ell_f, \ell_B \geq 0$, and functions $\varphi_1, \varphi_2, f_1, f_2 \geq 0$.

Lemma

Assume $1 > a > \delta \geq 0$, $\alpha \geq 2 - \delta$, and $|u|^\alpha \in W^{1,1}(U)$. Let $r > 0$.

(i) For any $\varepsilon > 0$ one has

$$\int_{\Gamma} |u|^{\alpha+r} d\sigma \leq \varepsilon \int_U |u|^{\alpha-2+\delta} |\nabla u|^{2-a} dx + c_1 \int_U |u|^{\alpha+r} dx \\ + (c_2(\alpha + r))^{\frac{2-a}{1-a}} \varepsilon^{-\frac{1}{1-a}} \int_U |u|^{\alpha + \frac{(2-a)r+a-\delta}{1-a}} dx.$$

(ii) If $\alpha > \frac{n(r+a-\delta)}{2-a}$, then for any $\varepsilon > 0$ one has

$$\int_U |u|^{\alpha+r} dx \leq \varepsilon \int_U |u|^{\alpha-2+\delta} |\nabla u|^{2-a} dx + D_1 \|u\|_{L^\alpha}^{\alpha+r} \\ + D_2 \varepsilon^{-\frac{\theta}{1-\theta}} \|u\|_{L^\alpha}^{\alpha \left(1 + \frac{2-a}{n} \cdot \frac{\theta}{1-\theta}\right)}.$$

Estimates of the Lebesgue norms

If α is large and $t > 0$ then

$$\begin{aligned} \frac{d}{dt} \int_U u(x, t)^\alpha dx + \int_U |\nabla u(x, t)|^{2-a} u(x, t)^{\alpha-\lambda-1} dx \\ \leq C_0 \cdot \left(\|u(t)\|_{L^\alpha(U)}^{\nu_1} + \|u(t)\|_{L^\alpha(U)}^{\nu_2} + \Upsilon(t) \right), \end{aligned}$$

where $\Upsilon(t) = \|\varphi_1(t)\|_{L^{q_1}(\Gamma)}^{q_1} + \|\varphi_2(t)\|_{L^{q_2}(\Gamma)}^{q_2} + \|f_1(t)\|_{L^{q_3}(U)}^{q_3} + \|f_2(t)\|_{L^{q_4}(U)}^{q_4}$.

Theorem (Celik-H.-Kieu 2018)

If $\int_0^T (1 + \Upsilon(t)) dt \leq C_1 \cdot (1 - 2^{-\nu_4}) \left(1 + \int_U u_0^\alpha(x) dx\right)^{-\nu_4}$, then

$$\int_U u^\alpha(x, t) dx \leq 1 + 2 \int_U u_0^\alpha(x) dx \quad \text{for all } t \in [0, T],$$

$$\int_0^T \int_U |\nabla u|^{2-a} u^{\alpha-\lambda-1} dx dt \leq 2^{\nu_4+1} (1 + 1/\nu_4) \left(1 + \int_U u_0^\alpha(x) dx\right).$$

Lemma

If α is large and $T > T_2 > T_1 \geq 0$ then

$$\begin{aligned} & \sup_{t \in [T_2, T]} \int_U u^\alpha(x, t) dx + \int_{T_2}^T \int_U |\nabla u(x, t)|^{2-a} u(x, t)^{\alpha-\lambda-1} dx dt \\ & \leq c_6(1+T) \left(1 + \frac{1}{T_2 - T_1}\right) \alpha^2 \mathcal{M}_0 \left(\|u\|_{L^{\tilde{\kappa}\alpha}(U \times (T_1, T))}^{\nu_5} + \|u\|_{L^{\tilde{\kappa}\alpha}(U \times (T_1, T))}^{\nu_6} \right), \end{aligned}$$

where $\mathcal{M}_0 = 1 + \|\varphi_1\|_{L^{q_1}(\Gamma_T)}^{\frac{2-a}{1-a}} + \|\varphi_2\|_{L^{q_2}(\Gamma_T)}^{\frac{2-a}{1-a}} + \|f_1\|_{L^{q_3}(Q_T)} + \|f_2\|_{L^{q_4}(Q_T)}$.

Parabolic Sobolev embedding

Lemma

Assume $1 > a > \delta \geq 0$, $\alpha > 0$ is large, and $T > 0$, then

$$\begin{aligned} & \left(\int_0^T \int_U |u|^{\kappa\alpha} dx dt \right)^{\frac{1}{\kappa\alpha}} \\ & \leq (c_5 \alpha^{2-a})^{\frac{1}{\kappa\alpha}} \left(\int_0^T \int_U |u|^{\alpha+\delta-a} dx dt + \int_0^T \int_U |u|^{\alpha+\delta-2} |\nabla u|^{2-a} dx dt \right)^{\frac{\tilde{\theta}}{\alpha+\delta-a}} \\ & \quad \cdot \sup_{t \in [0, T]} \left(\int_U |u(x, t)|^\alpha dx \right)^{\frac{1-\tilde{\theta}}{\alpha}}, \end{aligned}$$

where $c_5 \geq 1$ is independent of α and T , and

$$\tilde{\theta} = \tilde{\theta}_\alpha \stackrel{\text{def}}{=} \frac{1}{1 + \frac{\alpha(2-a)}{n(\alpha+\delta-a)}}, \quad \kappa = \kappa(\alpha) \stackrel{\text{def}}{=} 1 + \frac{2-a}{n} - \frac{a-\delta}{\alpha} = 1 + (a-\delta) \left(\frac{1}{\alpha_*} - \frac{1}{\alpha} \right)$$

Lemma

If $T > T_2 > T_1 \geq 0$ then

$$\|u\|_{L^{\kappa\alpha}(U \times (T_2, T))} \leq A_\alpha^{\frac{1}{\alpha}} \left(\|u\|_{L^{\tilde{\kappa}\alpha}(U \times (T_1, T))}^{\nu_5} + \|u\|_{L^{\tilde{\kappa}\alpha}(U \times (T_1, T))}^{\nu_7} \right)^{\frac{1}{\alpha}},$$

where $A_\alpha = c_7(1+T)^2 \left(1 + \frac{1}{T_2 - T_1}\right)^2 \alpha^{6-a} \mathcal{M}_0^2$.

Fix $\tilde{\kappa}$.

Set up $\kappa = \kappa(\alpha) > \tilde{\kappa}$.

Then we modify Moser's iterations.

Theorem (Celik-H.-Kieu 2018)

Let α_0 be sufficiently large. If $T > 0$ and $\sigma \in (0, 1)$ then

$$\|u\|_{L^\infty(U \times (\sigma T, T))} \leq C \left(1 + \frac{1}{\sigma T}\right)^{\omega_1} (1 + T)^{\omega_2} \mathcal{M}_0^{\omega_3} \\ \cdot \max \left\{ \|u\|_{L^{\tilde{\mu}}_{\tilde{\kappa}\alpha_0}(U \times (0, T))}, \|u\|_{L^{\tilde{\nu}}_{\tilde{\kappa}\alpha_0}(U \times (0, T))} \right\}.$$

Theorem (Celik-H.-Kieu 2018)

Let α_0 be sufficiently large. If $T > 0$ is small, then for $0 < \varepsilon < \min\{1, T\}$, one has

$$\|u\|_{L^\infty(U \times (\varepsilon, T))} \leq C \varepsilon^{-\omega_1} (1 + T)^{\omega_2 + \frac{\tilde{\nu}}{\beta_1}} (1 + \|u_0(x)\|_{L^{\beta_1}(U)})^{\tilde{\nu}} \mathcal{M}_0^{\omega_3}.$$

Assumption (A2).

- (i) The function $Z(u)$ satisfies

$$|Z'(u)| \leq d_4 u^{\ell_Z - 1} \quad \forall u \in (0, \infty),$$

for some constant $d_4 > 0$.

- (ii) There are non-negative functions $\varphi_3(x, t)$ and $\varphi_4(x, t)$ defined on $\Gamma \times (0, \infty)$ such that

$$\left| \frac{\partial B(x, t, u)}{\partial t} \right| \leq \varphi_3(x, t) + \varphi_4(x, t) u^{\ell_B} \quad \forall (x, t, u) \in \Gamma \times [0, \infty) \times [0, \infty).$$

- (iii) We also assume

$$|B(x, t, u)| \leq \varphi_1(x, t) + \varphi_2(x, t) u^{\ell_B} \quad \forall (x, t, u) \in \Gamma \times [0, \infty) \times [0, \infty),$$

$$|f(x, t, u)| \leq f_1(x, t) + f_2(x, t) u^{\ell_f} \quad \forall (x, t, u) \in U \times [0, \infty) \times [0, \infty).$$

Assumption (A3).

$$2\ell_Z > \lambda + 1.$$

Remark. For our original problem, $\ell_Z = 2\lambda$, then Assumption (A3) becomes $\lambda > 1/3$.

- For slightly compressible fluids, $\lambda = 1$.
- For ideal gases, $\lambda = 1/2$.
- For isentropic gas flows, all values of the specific heat ratio γ found belong to the interval $(1, 2)$, therefore $\lambda = 1/(1 + \gamma)$, satisfies

$$1/3 < \lambda < 1/2.$$

Thus Assumption (A3) is naturally met in all cases.

Let

$$\mathcal{I}(t) = \int_U H(|\nabla u(x, t) + Z(u(x, t))|) dx,$$

$$\mathcal{Z}_0 = \int_U u_0^{\lambda+1}(x) dx + \mathcal{I}(0) + \int_{\Gamma} (\varphi_1(x, 0)u_0(x) + \varphi_2(x, 0)u_0^{\ell_B+1}(x)) d\sigma,$$

$$N_1(t) = \int_{\Gamma} (\varphi_1^{\eta_5}(x, t) + \varphi_2^2(x, t)) d\sigma,$$

$$N_2(t) = N_1(t) + \int_{\Gamma} (\varphi_3^{\eta_5}(x, t) + \varphi_4^2(x, t)) d\sigma + \int_U (f_1^{\eta_6}(x, t) + f_2^4(x, t)) dx.$$

Theorem

If $T > 0$ is small, then for all $t \in (0, T]$

$$\begin{aligned} \int_U |\nabla u(x, t)|^{2-a} dx \leq C \left\{ \mathcal{Z}_0 + (t+1) \left(1 + \int_U u_0^{\eta_7}(x) dx \right) \right. \\ \left. + N_1(t) + \int_0^t N_2(\tau) d\tau \right\}. \end{aligned}$$

4. Flows of mixed regimes

- Unified models
- Estimates for solutions
- Continuous dependence on the boundary data
- Structural stability

- Darcy:

$$\mathbf{v} = -k\nabla p.$$

- Pre-Darcy: When $|\mathbf{v}|$ is small,

$$|\mathbf{v}|^{-\alpha}\mathbf{v} = -k\nabla p, \alpha \in (0, 1).$$

- Post-Darcy:

$$(a_0 + a_1|\mathbf{v}|^{\alpha_1} + \dots + a_N|\mathbf{v}|^{\alpha_N})\mathbf{v} = -\nabla p.$$

$$-\nabla p = \mathbf{G}(v) = \begin{cases} g(|v|)v & \text{if } v \in \mathbb{R}^n \setminus \{0\}, \\ 0 & \text{if } v = 0, \end{cases}$$

where $g(s) > 0$ on $(0, \infty)$ and $\lim_{s \searrow 0} sg(s) = 0$.

Solve for v . Taking the modulus both sides, we have $G(|v|) = |\nabla p|$, where

$$G(s) = \begin{cases} sg(s) & \text{if } s > 0, \\ 0 & \text{if } s = 0. \end{cases}$$

We assume

- $G(s)$ is strictly increasing on $[0, \infty)$,
- $G(s) \rightarrow \infty$ as $s \rightarrow \infty$, and
- the function $1/g(s)$ on $(0, \infty)$ can be extended to a continuous function $k_g(s)$ on $[0, \infty)$.

Then

$$v = -K(|\nabla p|)\nabla p,$$

where $K(\xi) = k_g(G^{-1}(\xi))$ for $\xi \geq 0$.

Two models of g

Model 1. Function $g(s)$ is piece-wise defined:

$$g(s) = \bar{g}(s) \stackrel{\text{def}}{=} c_1 s^{-\alpha} \mathbf{1}_{(0, s_1)}(s) + c_2 \mathbf{1}_{[s_1, s_2]}(s) + g_F(s) \mathbf{1}_{(s_2, \infty)}(s) \quad \text{for } s > 0,$$

Continuity condition:

$$c_1 s_1^{-\alpha} = c_2 = g_F(s_2).$$

$$K(\xi) = \bar{K}(\xi) \stackrel{\text{def}}{=} M_1 \xi^{\beta_1} \mathbf{1}_{[0, Z_1]}(\xi) + M_2 \mathbf{1}_{[Z_1, Z_2]}(\xi) + K_F(\xi) \mathbf{1}_{(Z_2, \infty)}(\xi).$$

Model 2. Function $g(s)$ is smooth on $(0, \infty)$:

$$g(s) = g_I(s) \stackrel{\text{def}}{=} a_{-1} s^{-\alpha} + a_0 + a_1 s^{\alpha_1} + \dots + a_N s^{\alpha_N} \quad \text{for } s > 0, \quad (4.1)$$

where $N \geq 1$, $\alpha \in (0, 1)$, $\alpha_N > 0$,

$$a_{-1}, a_N > 0 \text{ and } a_i \geq 0 \quad \forall i = 0, 1, \dots, N-1.$$

$$K(\xi) = K_I(\xi) \stackrel{\text{def}}{=} \frac{s(\xi)^\alpha}{a_{-1} + a_0 s(\xi)^\alpha + a_1 s(\xi)^{\alpha+\alpha_1} + \dots + a_N s(\xi)^{\alpha+\alpha_N}},$$

with $G(s(\xi)) = \xi$.

Two direct models of K

$$\mathbf{v} = -K(|\nabla p|)\nabla p.$$

Note: $K(\xi)$ behaves like ξ^{β_1} for small ξ , and like $(1 + \xi)^{-\beta_2}$ for large ξ ,

$$\beta_1 = \frac{\alpha}{1 - \alpha}, \quad \beta_2 = \frac{\alpha_N}{\alpha_N + 1}.$$

Model 3.

$$K(\xi) = \hat{K}(\xi) \stackrel{\text{def}}{=} \frac{a\xi^{\beta_1}}{(1 + b\xi^{\beta_1})(1 + c\xi^{\beta_2})}.$$

Model 4. More precisely, $K(\xi)$ is close to $M_1\xi^{\beta_1}$ when $\xi \rightarrow 0$, and to $K_F(\xi)$ when $\xi \rightarrow \infty$. Then we choose

$$K(\xi) = K_M(\xi) \stackrel{\text{def}}{=} K_F(\xi) \cdot \frac{\bar{k}\xi^{\beta_1}}{1 + \bar{k}\xi^{\beta_1}}.$$

where $\bar{k} = M_1/K_F(0) > 0$.

Initial Boundary Value Problem

Let $K(\xi)$ be one of the functions $\bar{K}(\xi)$, $K_I(\xi)$, $\hat{K}(\xi)$, $K_M(\xi)$.

Slightly compressible fluids (with simplification.)

After scaling the time variable (to simplify ϕ), we obtain the IBVP:

$$\begin{cases} p_t = \nabla \cdot (K(|\nabla p|)\nabla p) & \text{in } U \times (0, \infty), \\ p(x, 0) = p_0(x), & \text{in } U \\ p = \psi(x, t), & \text{on } \partial U \times (0, \infty). \end{cases}$$

Lemma (Degeneracy)

Then there exist $d_2, d_3 > 0$ such that

$$\frac{d_2 \xi^{\beta_1}}{(1 + \xi)^{\beta_1 + \beta_2}} \leq K(\xi) \leq \frac{d_3 \xi^{\beta_1}}{(1 + \xi)^{\beta_1 + \beta_2}} \quad \forall \xi \geq 0.$$

Energy estimates

Let $\bar{p} = p - \Psi$, where $\Psi(x, t)$ is an extension of ψ from $x \in \partial U$ to $x \in \bar{U}$.

Theorem (Celik-H.-Ibragimov-Kieu 2017)

(i) There exists a positive constant C such that for all $t \geq 0$,

$$\|\bar{p}(t)\|^2 \leq \|\bar{p}(0)\|^2 + C [1 + \text{Env}(f(t))]^{\frac{2}{2-\beta_2}},$$

where $f(t) = f[\Psi](t) \stackrel{\text{def}}{=} \|\nabla \Psi(t)\|^2 + \|\Psi_t(t)\|^{\frac{2-\beta_2}{1-\beta_2}}$.

(ii) Furthermore,

$$\limsup_{t \rightarrow \infty} \|\bar{p}(t)\|^2 \leq C(1 + \limsup_{t \rightarrow \infty} f(t))^{\frac{2}{2-\beta_2}}.$$

(iii) If $\lim_{t \rightarrow \infty} \|\nabla \Psi(t)\| = \lim_{t \rightarrow \infty} \|\Psi_t(t)\| = 0$, then

$$\lim_{t \rightarrow \infty} \|\bar{p}(t)\| = 0.$$

Theorem (Celik-H.-Ibragimov-Kieu 2017)

(i) For all $t \geq 0$,

$$\int_U |\nabla p(x, t)|^{2-\beta_2} dx \leq C \left(1 + \|\bar{p}(0)\|^2 + e^{-\frac{t}{2}} \int_U |\nabla p(x, 0)|^{2-\beta_2} dx + [Env(f(t))]^{\frac{2}{2-\beta_2}} + \int_0^t e^{-\frac{1}{2}(t-\tau)} \|\nabla \Psi_t(\tau)\|^2 d\tau \right).$$

(ii) $\limsup_{t \rightarrow \infty} \int_U |\nabla p(x, t)|^{2-\beta_2} dx \leq C(1 + \limsup_{t \rightarrow \infty} G_1(t))$,

where $G_1(t) = G_1[\Psi](t) \stackrel{\text{def}}{=} f(t)^{\frac{2}{2-\beta_2}} + \|\nabla \Psi_t(t)\|^2$.

(iii) If $\lim_{t \rightarrow \infty} \|\nabla \Psi(t)\| = \lim_{t \rightarrow \infty} \|\Psi_t(t)\| = \lim_{t \rightarrow \infty} \|\nabla \Psi_t(t)\| = 0$ then

$$\lim_{t \rightarrow \infty} \int_U |\nabla p(x, t)|^{2-\beta_2} dx = 0.$$

Theorem (Celik-H.-Ibragimov-Kieu 2017)

(i) If $t \geq 1$ then

$$\int_U |\nabla p(x, t)|^{2-\beta_2} dx \leq C \left(1 + \|\bar{p}(0)\|^2 + [Env(f(t))]^{\frac{2}{2-\beta_2}} + \int_{t-1}^t \|\nabla \Psi_t(\tau)\|^2 d\tau \right)$$

(ii) One has

$$\begin{aligned} \limsup_{t \rightarrow \infty} \int_U |\nabla p(x, t)|^{2-\beta_2} dx \\ \leq C \left(1 + \limsup_{t \rightarrow \infty} f(t)^{\frac{2}{2-\beta_2}} + \limsup_{t \rightarrow \infty} \int_{t-1}^t \|\nabla \Psi_t(\tau)\|^2 d\tau \right). \end{aligned}$$

(iii) Moreover, $\lim_{t \rightarrow \infty} \int_U |\nabla p(x, t)|^{2-\beta_2} dx = 0$ provided

$$\lim_{t \rightarrow \infty} \|\nabla \Psi(t)\| = \lim_{t \rightarrow \infty} \|\Psi_t(t)\| = \lim_{t \rightarrow \infty} \int_{t-1}^t \|\nabla \Psi_t(\tau)\|^2 d\tau = 0.$$

Continuous dependence on the boundary data

- We consider $K(\xi) = \bar{K}(\xi)$, $K_I(\xi)$, $\hat{K}(\xi)$, $K_M(\xi)$.
- For $i = 1, 2$, let $p_i(x, t)$ be a solution with boundary data $\psi_i(x, t)$, with extensions $\Psi_i(x, t)$. Let $\bar{p}_i = p_i - \Psi_i$.
- Denote $\Phi = \Psi_1 - \Psi_2$ and $\bar{P} = \bar{p}_1 - \bar{p}_2 = p_1 - p_2 - \Phi$.
- Set $\mathcal{Y}_0 = 1 + \sum_{i=1,2} \left(\|\bar{p}_i(0)\|^2 + \|\nabla p_i(0)\|_{L^{2-\beta_2}}^{2-\beta_2} \right)$,

$$\begin{aligned} \tilde{\mathcal{Y}}(t) &= \mathcal{Y}_0 + \sum_{i=1,2} [Env(f[\Psi_i](t))]^{\frac{2}{2-\beta_2}} \\ &+ \begin{cases} \int_0^t e^{-\frac{1}{2}(t-\tau)} \sum_{i=1,2} \|\nabla \Psi_{i,t}(\tau)\|^2 d\tau & \text{if } 0 \leq t < 1, \\ \int_{t-1}^t \sum_{i=1,2} \|\nabla \Psi_{i,t}(\tau)\|^2 d\tau & \text{if } t \geq 1. \end{cases} \end{aligned}$$

- Let

$$D(t) = \|\Phi_t(t)\| + \|\nabla \Phi(t)\|_{L^{2-\beta_2}} + \|\nabla \Phi(t)\|_{L^{2+\beta_1}}^{2+\beta_1}.$$

and $\mathcal{D} = \limsup_{t \rightarrow \infty} D(t)$.

- For asymptotic estimates, we use

$$\tilde{\mathcal{A}} = \left(\sum_{i=1,2} \limsup_{t \rightarrow \infty} f[\Psi_i](t) \right)^{\frac{1}{2-\beta_2}},$$

$$\tilde{\mathcal{K}} = \tilde{\mathcal{A}}^2 + \sum_{i=1,2} \limsup_{t \rightarrow \infty} \int_{t-1}^t \|\nabla \Psi_{i,t}(\tau)\|^2 d\tau.$$

Theorem (Celik-H.-Ibragimov-Kieu 2017)

For $t \geq 0$,

$$\|\bar{P}(t)\|^2 \leq \|\bar{P}(0)\|^2 + C \left\{ \text{Env} \left[\tilde{\mathcal{Y}}(t)^{\frac{\beta_1+\beta_2}{2-\beta_2} + \frac{1}{2}} D(t) \right] \right\}^{\frac{2}{2+\beta_1}}.$$

If $\tilde{\mathcal{K}} < \infty$ then

$$\limsup_{t \rightarrow \infty} \|\bar{P}(t)\|^2 \leq C \left\{ (1 + \tilde{\mathcal{K}})^{\frac{\beta_1+\beta_2}{2-\beta_2} + \frac{1}{2}} D \right\}^{\frac{2}{2+\beta_1}}.$$

Structural stability

- Consider $K(\xi) = K_I(\xi, \vec{a})$ and study the dependence of the solutions on the coefficient vector \vec{a} .
- Let $N \geq 1$ and the exponent vector $\vec{\alpha} = (-\alpha, 0, \alpha_1, \dots, \alpha_N)$ be fixed.
- Denote the set of admissible \vec{a}

$$S = \{\vec{a} = (a_{-1}, a_0, \dots, a_N) : a_{-1}, a_N > 0, a_0, a_1, \dots, a_{N-1} \geq 0\}.$$

Lemma (Perturbed Monotonicity)

For any coefficient vectors $\vec{a}^{(1)}, \vec{a}^{(2)} \in S$, and any $y, y' \in \mathbb{R}^n$, one has

$$\begin{aligned} (K_I(|y'|, \vec{a}^{(1)})y' - K_I(|y|, \vec{a}^{(2)})y) \cdot (y' - y) &\geq \frac{d_6 |y - y'|^{2+\beta_1}}{(1 + |y| + |y'|)^{\beta_1 + \beta_2}} \\ &\quad - d_7 K(|y| \vee |y'|, \vec{a}^{(1)} \wedge \vec{a}^{(2)}) (|y| \vee |y'|) |\vec{a}^{(1)} - \vec{a}^{(2)}| |y - y'|. \end{aligned}$$

$$\mathcal{Y}(t) = \mathcal{Y}_0 + [\text{Env}(f(t))]^{\frac{2}{2-\beta_2}} + \begin{cases} \int_0^t \|\nabla \Psi_t(\tau)\|^2 d\tau & \text{if } 0 \leq t < 1, \\ \int_{t-1}^t \|\nabla \Psi_t(\tau)\|^2 d\tau & \text{if } t \geq 1, \end{cases}$$

$$\mathcal{A} = \limsup_{t \rightarrow \infty} f(t)^{\frac{1}{2-\beta_2}} \quad \text{and} \quad \mathcal{K} = \mathcal{A}^2 + \limsup_{t \rightarrow \infty} \int_{t-1}^t \|\nabla \Psi_t(\tau)\|^2 d\tau.$$

Theorem (Celik-H.-Ibragimov-Kieu 2017)

(i) For $t \geq 0$, one has

$$\int_U |P(x, t)|^2 dx \leq \int_U |P(x, 0)|^2 dx + C[\text{Env}(\mathcal{Y}(t))]^{\frac{2}{2-\beta_2}} |\vec{a}^{(1)} - \vec{a}^{(2)}|^{\frac{2}{2+\beta_1}}.$$

(ii) If $\mathcal{K} < \infty$ then

$$\limsup_{t \rightarrow \infty} \int_U |P(x, t)|^2 dx \leq C(1 + \mathcal{K})^{\frac{2}{2-\beta_2}} |\vec{a}^{(1)} - \vec{a}^{(2)}|^{\frac{2}{2+\beta_1}}.$$

5. Publications and Funding

- 1 D. Cao, L. Hoang, *Asymptotic expansions in a general system of decaying functions for solutions of the Navier-Stokes equations*, 54 pp, submitted.
- 2 D. Cao, L. Hoang, *Long-time asymptotic expansions for Navier-Stokes equations with power-decaying forces*, Proceedings A of the Royal Society of Edinburgh, 35 pp, accepted.
- 3 L. Hoang, T. Kieu, *Global estimates for generalized Forchheimer flows of slightly compressible fluids*, Journal d'Analyse Mathématique, (2018), 54 pp, in press.
- 4 E. Celik, L. Hoang, T. Kieu, *Doubly nonlinear parabolic equations for a general class of Forchheimer gas flows in porous media*, Nonlinearity, Vol. 31, No. 8 (2018) 3617–3650.
- 5 C. Foias, L. Hoang, J.-C. Saut, *Navier and Stokes meet Poincaré and Dulac*, J. Appl. Anal. Comput., Volume 8, Number 3, (June 2018) 727–763. (survey)
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- 1 Nonlinear Couplings for Flows in Fractured Porous Media: Analysis and Numerical Algorithms**
National Science Foundation - Applied Mathematics (DMS 1412796)
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