# Minimal Harmonic Measure on Complementary Regions 

Roger W. Barnard, Leah Cole, and Alexander Yu. Solynin

(Communicated by Doron Lubinsky)


#### Abstract

For any two points $a_{1}$ and $a_{2}$ in an open disk $\Delta$ on the complex sphere $\overline{\mathbb{C}}$, let $L$ be a curve separating $a_{1}$ from $a_{2}$ on $\overline{\mathbb{C}}$, which splits $\overline{\mathbb{C}}$ into two complementary regions $B_{1} \ni a_{1}$ and $B_{2} \ni a_{2}$. Let $l$ be the part of this curve lying in $\bar{\Delta}$. In this note we study how small the average harmonic measure $$
\frac{1}{2}\left(\omega\left(a_{1}, l, B_{1}\right)+\omega\left(a_{2}, l, B_{2}\right)\right)
$$ can be. This question can be interpreted as a problem on the minimal average temperature at two points of a long cylinder composed of two media separated by a heating membrane each of which contains a reference point.


Keywords. Harmonic measure, module of a quadrilateral, complete elliptic integral.

2000 MSC. 30C85, 33E05.

## 1. Introduction to the problem

Let a closed (non-Jordan in general) curve $L$ split the complex sphere $\overline{\mathbb{C}}$ into two complementary regions $B_{1}$ and $B_{2}$ containing points $a_{1}$ and $a_{2}$ respectively. Let $\Delta$ be an open disk containing $a_{1}$ and $a_{2}$ and let $l_{1}, l_{2}$ be parts of boundaries $\partial B_{1}$ and $\partial B_{2}$ lying in $\bar{\Delta}$.
How small compared to $L$ can the set $l=l_{1} \cup l_{2}$ be? This is the main question studied in this note. For a measure of "smallness" it is natural to choose the average of harmonic measures of the sets $l_{1}$ and $l_{2}$ evaluated at the reference points $a_{1}$ and $a_{2}$. In this way we arrive at the following extremal problem.

Problem $\mathcal{P}_{\Delta}\left(a_{1}, a_{2}\right)$. Find the minimization of the sum

$$
\begin{equation*}
\frac{1}{2}\left(\omega\left(a_{1}, l_{1}, B_{1}\right)+\omega\left(a_{2}, l_{2}, B_{2}\right)\right) \rightarrow \inf . \tag{1.1}
\end{equation*}
$$

Received May 13, 2002.
The research of the third author was supported in part by Russian Fund for Fundamental Research, grant no. 00-01-00118a.

Here $\omega(z, E, D)$ denotes the harmonic measure of the set $E \subset \partial D$ with respect to the region $D$ evaluated at the point $z$. That is, $\omega(z, E, D)$ is the bounded harmonic function having boundary values 1 a.e. on $E$ and 0 a.e. on $\partial D \backslash E$. The infimum in (1.1) is taken among all curves $L$ separating $a_{1}$ from $a_{2}$ on $\overline{\mathbb{C}}$ as it is illustrated in Figure 1. For properties of the harmonic measure we refer to [A1, Ch. 1.1] or [G, Ch. 8].


Figure 1. Competing domains of Problem $\mathcal{P}_{\Delta}\left(a_{1}, a_{2}\right)$.
Putting more data one can additionally require that $L$ traces through a prescribed set of contact points $b_{1}, \ldots, b_{n}$ lying in $\Delta^{\prime}=\bar{\Delta} \backslash\left\{a_{1}, a_{2}\right\}$. In this paper we deal with the simplest case when the set of contact points is empty. A set of similar extremal problems on the harmonic measure was studied in [S1].
Since the harmonic measure is conformally invariant, the problem formulated above is invariant under Möbius maps. Therefore in what follows we may assume without loss of generality that $\Delta$ is the unit disk $\Delta=\{z:|z|<1\}$ and $a_{1}=r$, $a_{2}=-r$, where $r, 0<r<1$ is defined by

$$
\log \frac{1+r}{1-r}=\rho_{\Delta}\left(a_{1}, a_{2}\right)
$$

where $\rho_{\Delta}\left(a_{1}, a_{2}\right)$ stands for the hyperbolic distance between $a_{1}$ and $a_{2}$ with respect to the disk $\Delta$, see [A1, G]. Thus the Problem $\mathcal{P}_{\Delta}\left(a_{1}, a_{2}\right)$ actually depends on one positive parameter $0<r<1$ only.
In Section 2, we use results from the theory of modules of families of curves to give a qualitative solution of Problem $\mathcal{P}_{\Delta}(a, b)$ by including the minimal harmonic measure into a one-parameter family of functions. Then, in Section 3 we show that for every value of the parameter $r, 0<r<1$, except for one special value $r_{0}=g-\sqrt{g}=.34601 \ldots$, where $g=(1+\sqrt{5}) / 2$ is the golden ratio, the corresponding one-parametric family of configurations contains a unique minimal configuration that is either symmetric or degenerate. In the last section we consider possible generalizations and discuss some open questions.

## 2. Problem $\mathcal{P}(r)$

2.1. Reduction of the problem. First we reduce the problem considered to a simpler one. Note that the assumptions of the problem imply the existence of a continuum (i.e. connected compact set) $E \subset \bar{\Delta} \cap L$, which splits $\Delta$ into two disjoint regions $D_{1} \ni r$ and $D_{2} \ni-r$. Thus, $\Delta \backslash E=D_{1} \cup D_{2}$. By the Carleman Principle for the harmonic measure (or essentially by the Maximum Principle for harmonic functions),

$$
\frac{1}{2}\left(\omega\left(r, l_{1}, B_{1}\right)+\omega\left(-r, l_{2}, B_{2}\right)\right) \geq \frac{1}{2}\left(\omega\left(r, E, D_{1}\right)+\omega\left(-r, E, D_{2}\right)\right)
$$

and therefore the infimum of Problem $\mathcal{P}_{\Delta}\left(a_{1}, a_{2}\right)$ (1.1) is greater or equal to the infimum

$$
\begin{equation*}
\Omega(r)=\inf \frac{1}{2}\left(\omega\left(r, E, D_{1}\right)+\omega\left(-r, E, D_{2}\right)\right) \tag{2.1}
\end{equation*}
$$

taken among the set of all continua $E \subset \bar{\Delta}$ that split $\Delta$ into a pair of regions described above. Standard approximation arguments applied to the extremal configurations of Problem (2.1) described in Theorem 1, Section 3 show that the infima in (1.1) and (2.1) actually coincide. In what follows the minimization Problem (2.1) will be denoted by $\mathcal{P}(r)$.
One advantage of formulation (2.1) compared to (1.1) is that the infimum in (2.1) is realized for certain continuum $E^{*}$ or equivalently for a certain pair of non-overlapping regions $D_{1}^{*}, D_{2}^{*}$ one of which but not both may be degenerate.
2.2. Physical motivation. It might also be interesting to note that in formulation (2.1) our problem models a steady-state heat flow in a long cylindrical solid bar (one can have in mind, for instance, the plutonium bars in a nuclear reactor). Assume that the bar is composed of two media separated by a thin membrane having the same shape in each horizontal section as it is shown in Figure 2. Assume that we cool this bar along the cylindrical surface keeping the temperature at least $t_{0}$ at each of its points and assume that, in order to be operational, the membrane requires a temperature of at least $t_{1}>t_{0}$ at any of its points.
Assume next that in controlling the temperature at two points, $a_{1}$ in the first medium and $a_{2}$ in the second medium, we are interested in keeping the average temperature at these points as small as possible. What is the optimal design of the membrane that corresponds to the minimal average temperature? This question is equivalent to our Problem (2.1).


Figure 2. The minimal average temperature problem.
2.3. Extremal partition and module of a triad. Problem (2.1) is also connected with a problem on the extremal partition of $\Delta$ into a pair of disjoint quadrilaterals. Following J. Jenkins [J3], by a triad $(a, \alpha, D)$ we shall mean a simply connected region $D$ on $\overline{\mathbb{C}}$ with a fixed closed arc $\alpha$ on its boundary and a fixed point $a$ inside $D$. Let $\{Q\}$ be the set of all quadrilaterals $Q$ in $D^{\prime}=D \backslash\{a\}$ each of which has a distinguished pair of sides on the complementary boundary $\operatorname{arc} \beta=\partial D \backslash \alpha$ such that curves joining the distinguished sides of $Q$ separate $a$ from $\alpha$ inside $D$. Let $\bmod Q$ denote the module of $Q$ with respect to the family of curves in $D$ joining the distinguished sides of $Q$. Let $w=f(z)$ map $Q$ conformally onto a rectangle $\{w:|\operatorname{Re} w|<a,|\operatorname{Im} w|<b\}$ such that the distinguished sides of $Q$ map onto the vertical sides of the rectangle. Then, $\bmod Q=b / a$. For properties of the module of a quadrilateral we refer to [A1, A2, D, J2, S2].
It is well known [J1, J3], that there is a unique quadrilateral $Q_{0}=Q_{0}(a, \alpha, D)$ in the family $\{Q\}$ having the maximal module:

$$
\bmod Q_{0}=\max _{\{Q\}}\{\bmod Q\}
$$

Let $s$ be a conformal mapping of $D$ onto $\Delta$ normalized by the condition $s(a)=0$, $s(\alpha)=\left\{e^{i \theta}:|\theta-\pi| \leq \psi_{0}\right\}$ with $\psi_{0}=\pi \omega$, where $\omega=\omega(a, \alpha, D)$. Let $t=s^{-1}$. Then the extremal quadrilateral $Q_{0}$ is the $t$ - image of the quadrilateral $Q\left(\psi_{0}\right)=$ $\Delta \backslash[0,1]$ with vertices at the points $1+i 0,-e^{-i \psi_{0}},-e^{i \psi_{0}}$, and $1-i 0$.
By the module of a triad $(a, \alpha, D)$ we shall mean the module

$$
m\left(Q_{0}\right)=m\left(Q_{0}(a, \alpha, D)\right)
$$

The latter module and the harmonic measure $\omega=\omega(a, \alpha, D)$ are linked via a well known formula attributed to J. Hersch:

$$
\begin{equation*}
m\left(Q_{0}\right)=\frac{\mathbf{K}\left(\tan ^{2}(\pi(1-\omega) / 4)\right)}{\mathbf{K}^{\prime}\left(\tan ^{2}(\pi(1-\omega) / 4)\right)} \tag{2.2}
\end{equation*}
$$

where $\mathbf{K}$ denotes the complete elliptic integral of the first kind and

$$
\mathbf{K}^{\prime}(\tau)=\mathbf{K}\left(\sqrt{1-\tau^{2}}\right)
$$

Later on similar notation $\mathbf{E}$ and $\mathbf{E}^{\prime}$ will be used for the complete elliptic integrals of the second kind. In what follows it will be convenient to employ the notation

$$
\begin{equation*}
\mu(s)=\frac{\pi}{2} \frac{\mathbf{K}^{\prime}(s)}{\mathbf{K}(s)} \tag{2.3}
\end{equation*}
$$

for the quotient of the complete elliptic integrals, see [AVV, Chapter 5], which contains numerous useful properties of $\mu(s)$ and is accompanied by a software package useful in numerical computations concerning $\mu(s)$. Equation (2.2) shows that $m=m\left(Q_{0}\right)$ decreases from $+\infty$ to 0 when $\omega$ runs from 0 to 1 .

### 2.4. Harmonic measure and extremal quadrilaterals. Let

$$
\omega_{*}(r)=\omega(-r,[r, 1], \Delta) .
$$

Carleman's Principle and a symmetrization result of J. Krzyz [Kr], see also [B, Theorem 6.1], imply that

$$
\omega\left(-r, E, D_{2}\right) \geq \omega_{*}(r)
$$

with the sign of equality only for $E=[r, 1]$. For a fixed $\omega_{0}, \omega_{*}(r)<\omega_{0}<1$, consider Problem (2.1) under an additional constraint

$$
\begin{equation*}
\omega\left(-r, E, D_{2}\right)=\omega_{0} \tag{2.4}
\end{equation*}
$$

Our previous analysis shows that Problem (2.1), (2.4) is equivalent to the following problem on the modules. Let $Q_{1}, Q_{2}$ be a pair of disjoint quadrilaterals in $\Delta^{\prime}=\Delta \backslash\{ \pm r\}$ each of which has a pair of distinguished sides on $\mathbb{T}$. We assume in addition that curves $\gamma \subset Q_{1}$ joining the distinguished sides of $Q_{1}$ separate the point $r$ from $Q_{2}$ and the point $-r$ inside $\Delta^{\prime}$ and that similar curves $\gamma \subset Q_{2}$ separate the point $-r$ from $Q_{1}$ and $r$ inside $\Delta^{\prime}$. The problem is to find

$$
\bmod Q_{1} \rightarrow \max
$$

among all pairs of quadrilaterals $Q_{1}, Q_{2}$ satisfying the above mentioned conditions and such that

$$
\bmod Q_{2}=m_{0}=\frac{\pi}{2} \mu\left(\tan ^{2} \frac{\pi\left(1-\omega_{0}\right)}{4}\right)
$$

Let $\gamma_{1}$ denote the non-distinguished side of $Q_{1}$ that does not separate $Q_{1}$ from $Q_{2}$ and let $\gamma_{2}$ be a similar side of $Q_{2}$. Then we can consider a quadrilateral $Q_{1,2}=$ $\Delta \backslash\left(\gamma_{1} \cup \gamma_{2}\right)$ with $\gamma_{1}$ and $\gamma_{2}$ as its non-distinguished sides.
The quadrilateral $Q_{1,2}$ separates the points $r$ and $-r$ in $\Delta^{\prime}$. Therefore as a very special case of the Jenkins Theorem in [J1] we get the inequality

$$
\begin{equation*}
\bmod Q_{1,2} \leq \bmod Q(r) \tag{2.5}
\end{equation*}
$$

where $Q(r)$ denotes the quadrilateral $\Delta \backslash([-1,-r] \cup[r, 1])$ with the distinguished sides $\mathbb{T}^{+}=\{z \in \mathbb{T}: \operatorname{Im} z>0\}$ and $\mathbb{T}^{-}=\{z \in \mathbb{T}: \operatorname{Im} z<0\}$. Equality occurs in (2.5) if and only if $Q_{1,2}=Q(r)$. Inequality (2.5) follows also from a symmetrization result of O. Teichmüller, see [A2].
By Grötzsch's Lemma, see [J2, Ch.2],

$$
\bmod Q_{1}+\bmod Q_{2} \leq \bmod Q_{1,2}
$$

which together with 2.5 leads to the inequality

$$
\begin{equation*}
\bmod Q_{1} \leq \bmod Q(r)-m_{0} \tag{2.6}
\end{equation*}
$$

Let $f_{r}$ map $Q_{1,2}$ conformally onto a suitable rectangle

$$
\Pi(a)=\{w:|\operatorname{Im} w|<1,|\operatorname{Re} w|<a\}, \quad a>0,
$$



Figure 3. Extremal rectangles for $r=0.5$.
such that the slits $(-1,-r]$ and $[r, 1)$ correspond to the left and right vertical sides of $\Pi(a)$, respectively, and let $F_{r}$ be a function corresponding in a similar way to the extremal quadrilateral $Q(r)$. Equality in (2.6) occurs if and only if $Q_{1,2}=Q(r)$ and the sets $\Pi_{1}=f_{r}\left(Q_{1}\right)$ and $\Pi_{2}=f_{r}\left(Q_{2}\right)$ are two complementary rectangles in $\Pi(a)$, see Figure 3.
Combining all of these results we obtain Lemma 1, solving Problem (2.1), (2.4).
Lemma 1. For a fixed $\omega_{0}, \omega_{*}(r)<\omega_{0}<1$, there is a unique pair of domains $D_{1}^{*}, D_{2}^{*}$ minimizing the sum in (2.1) under the additional constraint

$$
\omega\left(-r, E, D_{2}\right)=\omega_{0}
$$

Moreover,

$$
\begin{aligned}
& D_{1}^{*}=D_{1}(a, b)=G_{r}(\Pi(b, a)) \cup[r, 1), \\
& D_{2}^{*}=D_{2}(a, b)=G_{r}(\Pi(-a, b)) \cup(-1,-r],
\end{aligned}
$$

where $G_{r}=F_{r}^{-1}$ and the parameters a and $b,-a<b<a<\infty$, are defined by the equations

$$
\begin{equation*}
a=\frac{\pi}{\mu\left(r^{2}\right)}, \quad b=\frac{\pi}{\mu\left(\tan ^{2} \frac{\pi\left(1-\omega_{0}\right)}{4}\right)}-\frac{\pi}{\mu\left(r^{2}\right)} . \tag{2.7}
\end{equation*}
$$

The minimum in (2.1) is

$$
\begin{equation*}
\Omega(a, b)=1-\frac{2}{\pi}\left(\tan ^{-1} \sqrt{\tau_{1}}+\tan ^{-1} \sqrt{\tau_{2}}\right) \tag{2.8}
\end{equation*}
$$

with $\tau_{1}$ and $\tau_{2}$ defined by

$$
\begin{equation*}
\frac{\pi}{\mu\left(\tau_{1}\right)}=a-b, \quad \frac{\pi}{\mu\left(\tau_{2}\right)}=a+b \tag{2.9}
\end{equation*}
$$

where $\mu(s)$ is defined in (2.3).

## 3. Minimal configurations

We have shown in Section 2 that the extremal configurations of Problem (2.1) are among the configurations $D_{1}(a, b), D_{2}(a, b)$ described by Lemma 1 with $a$ defined by (2.7) and some $b$ such that $-a \leq b \leq a$. Now we will study how the minimal average harmonic measure $\Omega(a, b)$ depends on the parameter $b$. Equations (2.8) and (2.9) show that for a fixed $a, \Omega(a, b)$ is an even function of $b$. Therefore the configuration $D_{1}(a,-b), D_{2}(a,-b)$ provides the extremal value for $\Omega(a, b)$ if and only if the configuration $D_{1}(a, b), D_{2}(a, b)$ does. This allows us to deal in what follows with non negative values of $b$ only. First we examine the limit cases.
Since $\Omega(a, b)$ is even in $b, b=0$ is always a critical point of $\Omega(a, b)$ and corresponds to the symmetric configuration $D_{1}(a, 0)=\Delta^{+}, D_{2}(a, 0)=\Delta^{-}$, where

$$
\Delta^{+}=\{z \in \Delta: \operatorname{Re} z>0\}, \quad \Delta^{-}=\Delta \backslash \bar{\Delta}^{+}
$$

Let $u_{s}(r)=\omega\left(r, I, D_{1}(a, 0)\right)$, where $I=[-i, i]$.
To find $\omega\left(r, I, D_{1}(a, 0)\right)$, we note that the function $\varphi(z)=i z\left(1-z^{2}\right)^{-1}$ maps $\Delta^{+}$ conformally onto the upper half-plane $\mathbb{H}$ such that $\varphi(I)=[-1 / 2,1 / 2]$ and $\varphi(r)=$ $\operatorname{ir}\left(1-r^{2}\right)^{-1}$. Since harmonic measure is invariant under conformal mapping,

$$
\omega\left(r, I, D_{1}(a, 0)\right)=\omega(\varphi(r),[-1 / 2,1 / 2], \mathbb{H}) .
$$

Let $2 \alpha \pi$ denote the angle formed by the segments $[-1 / 2, \varphi(r)]$ and $[1 / 2, \varphi(r)]$ at the point $\varphi(r)$. Then $2 \alpha=(2 / \pi) \tan ^{-1}\left(\left(1-r^{2}\right) / 2 r\right)$. It is well known [A1] that $\omega(\varphi(r),[-1 / 2,1 / 2], \mathbb{H})=2 \alpha$. Therefore,

$$
\begin{equation*}
u_{s}(r)=\frac{2}{\pi} \tan ^{-1} \frac{1-r^{2}}{2 r}=1-\frac{4}{\pi} \tan ^{-1} r . \tag{3.1}
\end{equation*}
$$

In the second limit case, $b=a$. In this case $D_{1}(a, a)=\emptyset, D_{2}(a, a)=\Delta(r)$, where $\Delta(r)=\Delta \backslash I(r), I(r)=[r, 1]$, where $a$ and $r$ are related via (2.7). Using a conformal mapping from $\Delta(r)$ onto $\mathbb{H}$,

$$
\varphi_{1}(z)=\varphi\left(\left(\frac{r-z}{1-r z}\right)^{1 / 2}\right)
$$

with $\varphi$ defined above, we can compute as in the symmetric case,

$$
\begin{equation*}
u_{d}(r)=\frac{1}{2}\left(1+\omega\left(r, I(r), D_{2}(a, a)\right)\right)=\frac{1}{2}+\frac{1}{\pi} \tan ^{-1} \frac{(1-r)^{2}}{2 \sqrt{2 r\left(1+r^{2}\right)}} \tag{3.2}
\end{equation*}
$$

The graphs of $u_{s}(r), u_{d}(r)$, and the graph of the function $u_{n}(r)$, which corresponds to a possible nondegenerate nonsymmetric configuration, as it will be explained later, are plotted in Figure 4.


Figure 4. Graphs of $u_{s}, u_{d}$, and $u_{n}$.

Lemma 2. For $0<r<1$, the equation $u_{s}(r)=u_{d}(r)$ has a unique solution $r_{0}=g-\sqrt{g}=.34601 \ldots$, where $g=(1+\sqrt{5}) / 2$ is the golden ratio and

$$
\begin{array}{lll}
u_{d}(r)<u_{s}(r) & \text { for } & 0<r<r_{0} \\
u_{d}(r)>u_{s}(r) & \text { for } & r_{0}<r<1
\end{array}
$$

Proof. By (3.1), (3.2), the equation $u_{s}(r)=u_{d}(r)$ is equivalent to

$$
\begin{equation*}
\tan ^{-1} \frac{1-r^{2}}{2 r}=\frac{\pi}{4}+\frac{1}{2} \tan ^{-1} \frac{(r-1)^{2}}{2 \sqrt{2 r\left(r^{2}+1\right)}} \tag{3.3}
\end{equation*}
$$

Taking the tangent of the both sides in (3.3), we obtain

$$
\frac{1-r^{2}}{2 r}=\frac{2 r^{2}+2-2 \sqrt{2 r^{3}+2 r}}{-4 r+2 \sqrt{2 r^{3}+2 r}}
$$

After some algebra this leads to the equation

$$
\begin{equation*}
r^{4}-2 r^{3}-2 r^{2}-2 r+1=0 \tag{3.4}
\end{equation*}
$$

Since

$$
\frac{d}{d r}\left(r^{4}-2 r^{3}-2 r^{2}-2 r+1\right)=2\left[2 r\left(r^{2}-1\right)-\left(3 r^{2}+1\right)\right]
$$

is clearly negative on $0<r<1$, (3.4) and therefore the equation $u_{s}(r)=u_{d}(r)$ has a unique solution $r_{0}$ in $0<r<1$. Solving equation (3.4) by radicals we obtain

$$
r_{0}=\frac{1+\sqrt{5}}{2}-\sqrt{\frac{1+\sqrt{5}}{2}}=.34601 \ldots
$$

This solution of (3.4) can be found also with Mathematica.
Lemma 2 shows, in particular, that the symmetric configuration is not extremal for $0<r<r_{0}$ and the degenerate configuration is not extremal for $r_{0}<r<1$.
To find values $\tau_{1}$ and $\tau_{2}$ minimizing the function (2.8), we consider a minimization problem

$$
\begin{equation*}
m\left(x_{1}, x_{2}\right)=1-\frac{2}{\pi}\left(\tan ^{-1} \sqrt{x_{1}}+\tan ^{-1} \sqrt{x_{2}}\right) \rightarrow \min \tag{3.5}
\end{equation*}
$$

under the constraint

$$
\begin{equation*}
\frac{1}{\mu\left(x_{1}\right)}+\frac{1}{\mu\left(x_{2}\right)}=\frac{2}{\mu\left(r^{2}\right)} \tag{3.6}
\end{equation*}
$$

Differentiating the Lagrange function

$$
\Phi\left(x_{1}, x_{2}\right)=1-\frac{2}{\pi}\left(\tan ^{-1} \sqrt{x_{1}}+\tan ^{-1} \sqrt{x_{2}}\right)-\lambda\left(\frac{1}{\mu\left(x_{1}\right)}+\frac{1}{\mu\left(x_{2}\right)}\right)
$$

of problem (3.5), (3.6) and using the differentiation formula

$$
\begin{equation*}
\mu^{\prime}(s)=\frac{-\pi^{2}}{4 s\left(1-s^{2}\right) \mathbf{K}^{\prime 2}(s)} \tag{3.7}
\end{equation*}
$$

see [AVV, p. 82], we find

$$
\begin{aligned}
-\frac{\pi}{2} \Phi_{x_{1}} & =\frac{1}{2 \sqrt{x_{1}}\left(1+x_{1}\right)}+\lambda \frac{\pi}{2 x_{1}\left(1-x_{1}^{2}\right) \mathbf{K}^{\prime 2}\left(x_{1}\right)} \\
-\frac{\pi}{2} \Phi_{x_{2}} & =\frac{1}{2 \sqrt{x_{2}}\left(1+x_{2}\right)}+\lambda \frac{\pi}{2 x_{2}\left(1-x_{2}^{2}\right) \mathbf{K}^{\prime 2}\left(x_{2}\right)}
\end{aligned}
$$

Therefore the critical values of $x_{1}, x_{2}$ of problem (3.5), (3.6) can be found from the equation

$$
\begin{equation*}
\sqrt{x}(1-x) \mathbf{K}^{\prime 2}(x)=t \tag{3.8}
\end{equation*}
$$

with $t=-\pi \lambda$.
Lemma 3. Let $f(x)=\sqrt{x}(1-x) \mathbf{K}^{\prime 2}(x)$. There is $x_{0}=.04771 \ldots$ such that $f(x)$ increases from 0 to $t_{0}=f\left(x_{0}\right)=4.08366 \ldots$ on $0<x<x_{0}$ and decreases from $t_{0}$ to 0 on $x_{0}<x<1$.

Proof. Let $g=\sqrt{f}$. Since

$$
\frac{d}{d x} \mathbf{K}(x)=\frac{\mathbf{E}(x)-\left(1-x^{2}\right) \mathbf{K}(x)}{x\left(1-x^{2}\right)},
$$

see [AVV, p. 50], differentiating we obtain

$$
g^{\prime}(x)=\frac{u(x)}{v(x)}
$$

where

$$
\begin{equation*}
u(x)=(1-x)^{2} \mathbf{K}^{\prime}(x)-4 \mathbf{E}^{\prime}(x), \quad v(x)=4 x^{3 / 4}(1+x) \sqrt{1-x} . \tag{3.9}
\end{equation*}
$$

Since $v(x)$ is positive on $0<x<1$, we need to show that $u(x)$ changes its sign only once in the interval $0<x<1$.
First we show that $u(x)$ is convex on $0<x<1$. Differentiating twice we find

$$
u^{\prime \prime}(x)=\frac{u_{1}(x)}{v_{1}(x)}
$$

where

$$
\begin{equation*}
u_{1}(x)=q(x) \mathbf{E}^{\prime}(x)\left(\frac{\mathbf{K}^{\prime}(x)}{\mathbf{E}^{\prime}(x)}+\frac{p(x)}{q(x)}\right) \tag{3.10}
\end{equation*}
$$

with

$$
\begin{aligned}
& p(x)=\frac{1}{4}(1+x)\left(1-x^{2}\right)\left(1+2 x-5 x^{2}\right), \\
& q(x)=\frac{1}{2} x^{2}(1+x)\left(1-x^{2}\right)\left(1-x+x^{2}\right),
\end{aligned}
$$

and

$$
v_{1}(x)=x^{2}\left(1-x^{2}\right)^{2} .
$$

Since $v_{1}(x)>0$ for $0<x<1$, the convexity of $u$ will follow if we show that $u_{1}(x)>0$ for $0<x<1$. Since $q(x)>0$ and $\mathbf{E}^{\prime}(x)>0$ for $0<x<1$, it follows from (3.10) that $u_{1}(x)>0$ if and only if,

$$
\begin{equation*}
\frac{\mathbf{K}^{\prime}(x)}{\mathbf{E}^{\prime}(x)}+\frac{p(x)}{q(x)}>0 \tag{3.11}
\end{equation*}
$$

for $0<x<1$. To prove (3.11), we apply the inequality

$$
\begin{equation*}
\frac{\mathbf{K}^{\prime}(x)}{\mathbf{E}^{\prime}(x)}>\frac{2}{1+x^{2}} \tag{3.12}
\end{equation*}
$$

which is equivalent to the known inequality

$$
\mathbf{K}(x)-\mathbf{E}(x)>\mathbf{E}(x)-\left(1-x^{2}\right) \mathbf{K}(x),
$$

see [AVV, Theorem 3.21(6)]. Thus, (3.11) follows from the inequality

$$
\frac{2}{1+x^{2}}+\frac{p(x)}{q(x)}=\frac{2+2 x-2 x^{3}-x^{4}}{x^{2}\left(1+x^{2}\right)\left(1-x+x^{2}\right)}>0
$$

which obviously holds for all $0<x<1$. Therefore $u_{1}(x)>0$ and $u(x)$ is convex on $0<x<1$.
Using the well known limit values of $\mathbf{K}^{\prime}(x)$ and $\mathbf{E}^{\prime}(x)$, see [AVV, p. 49], we find

$$
\lim _{x \rightarrow+0} u(x)=+\infty, \quad u(1)=-2 \pi
$$

which combined with the convexity property of $u(x)$ implies the desired monotonicity of $f$.
Solving the equation $u(x)=0$ with Maple, we find $x_{0}=.04771 \ldots$ and $f\left(x_{0}\right)=$ $t_{0}=4.08366 \ldots$

According to Lemma 3, for $0<t \leq t_{0}$, equation (3.8) has two solutions. Let $x=x_{1}(t)$ and $x=x_{2}(t)$, where $0<x_{1}(t)<x_{0}<x_{2}(t)<1$ for $t \neq t_{0}$. The graphs of $x_{1}, x_{2}$, and $f$, as defined in Lemma 3, are shown in Figure 5 .


Figure 5. Graphs of $x_{1}, x_{2}$, and $f$.

The analysis made above leads to the following three possibilities for the extremal configurations of Problem (2.1):

1. Extremal configuration is symmetric. In this case the minimal harmonic measure is defined by (3.1).
2. Extremal configuration is degenerate. In this case the minimal harmonic measure equals $u_{d}(r)$, where $u_{d}(r)$ is defined by (3.2).
3. In the third possible case the minimal harmonic measure is $m\left(\tau_{1}, \tau_{2}\right)$ with $m(\cdot, \cdot)$ defined by (3.5) and $\tau_{1}=x_{1}(t), \tau_{2}=x_{2}(t)$ for some $t, 0<t<t_{0}$, such that

$$
\begin{equation*}
\frac{1}{\mu\left(x_{1}(t)\right.}+\frac{1}{\mu\left(x_{2}(t)\right.}=\frac{2}{\mu\left(r^{2}\right)} . \tag{3.13}
\end{equation*}
$$

Let $G(t)$ denote the left hand side of (3.13). The graphs of $G(t)$ and $2 / \mu\left(r^{2}\right)$ are shown in Figure 6.
Now we prove our main result.


Figure 6. Graphs of $G(t)$ and $2 / \mu\left(r^{2}\right)$.

Theorem 1. Let $\Omega(r)$ denote the minimal harmonic measure of Problem $\mathcal{P}(r)$. Then $\Omega(r)=u_{d}(r)$ for $0<r \leq r_{0}$ and $\Omega(r)=u_{s}(r)$ for $r_{0} \leq r<1$.
For $r \neq r_{0}$, the extremal configuration is unique up to symmetry with respect to the imaginary axis. It is degenerate for $0<r<r_{0}$ and symmetric for $r_{0}<r<1$. For $r=r_{0}$, the degenerate and symmetric configurations are the only extremal configurations of Problem $\mathcal{P}(r)$.

Proof. Part 1: First we show that there is no nonsymmetric nondegenerate minimal configuration with the corresponding value of $t$ in the interval $\hat{t} \leq t \leq t_{0}$, where $\hat{t}=3.85609 \ldots$ is a solution to the equation

$$
\begin{equation*}
1-\frac{2}{\pi}\left(\tan ^{-1} \sqrt{x_{0}}+\tan ^{-1} \sqrt{x_{2}(t)}\right)=u_{d}(r) \tag{3.14}
\end{equation*}
$$

with $r=r(t)$ defined by

$$
\begin{equation*}
\frac{1}{\mu\left(x_{1}(t)\right.}+\frac{1}{\mu\left(x_{0}\right)}=\frac{2}{\mu\left(r^{2}\right)} . \tag{3.15}
\end{equation*}
$$

Since $\mu^{\prime}(s)<0$, see (3.7), the function $1 / \mu(s)$ is strictly increasing. Therefore $r(t)$ is an increasing function uniquely determined by (3.15). Now since the left hand side of (3.14) strictly increases in $t$ and the right hand side strictly decreases, $\hat{t}$ is uniquely determined.
Assume that for some $0<r^{\prime}<1$ there is a nonsymmetric nondegenerate minimal configuration, which corresponds to $t^{\prime}, \hat{t} \leq t^{\prime}<t_{0}$. Let $x_{1}\left(t^{\prime}\right)$ and $x_{2}\left(t^{\prime}\right)$ be the corresponding solutions of (3.8). Since these solutions are monotone, we have

$$
\begin{gather*}
x_{1}(\hat{t}) \leq x_{1}\left(t^{\prime}\right)<x_{1}\left(t_{0}\right)<x_{2}\left(t^{\prime}\right) \leq x_{2}(\hat{t}), \\
x_{1}(\hat{t})=.01836 \ldots, \quad x_{1}\left(t_{0}\right)=.04772 \ldots, \quad x_{2}(\hat{t})=.10536 \ldots . \tag{3.16}
\end{gather*}
$$

Therefore, it follows from (3.13) and (3.16) that

$$
r^{\prime}>r_{-}, \quad \text { where } r_{-}=r(\hat{t})=.17610 \ldots
$$

The corresponding average harmonic measure is defined by (2.8):

$$
\Omega\left(r^{\prime}\right)=1-\frac{2}{\pi}\left(\tan ^{-1} \sqrt{x_{1}\left(t^{\prime}\right)}+\tan ^{-1} \sqrt{x_{2}\left(t^{\prime}\right)}\right) .
$$

This combined with (3.16) leads to the lower bound

$$
\Omega\left(r^{\prime}\right)>1-\frac{2}{\pi}\left(\tan ^{-1} \sqrt{x_{0}}+\tan ^{-1} \sqrt{x_{2}(\hat{t})}\right)=u_{d}\left(r_{-}\right)>u_{d}\left(r^{\prime}\right)
$$

contradicting the conjectured extremality of the corresponding configuration.
Part 2: Next we will show that the function $G(t)$ strictly decreases from $+\infty$ to $G_{*}=G\left(t_{*}\right)=.45204 \ldots$, when $t$ runs from 0 to $t_{*}$, where $t_{*}=4.07389 \ldots$ is defined by the equation

$$
\frac{1}{\mu\left(x_{0}\right)}+\frac{1}{\mu\left(x_{2}\left(t_{*}\right)\right)}=G(\hat{t})
$$

Since $x_{1}(t)$ and $x_{2}(t)$ are solutions to (3.8), implicit differentiation gives

$$
G^{\prime}(t)=\frac{\pi}{t} \frac{f_{1}\left(x_{1}\right)+f_{1}\left(x_{2}\right)}{f_{1}\left(x_{1}\right) f_{1}\left(x_{2}\right)},
$$

where $x_{1}=x_{1}(t), x_{2}=x_{2}(t)$ and

$$
f_{1}(x)=2 \sqrt{x}(1+x) f^{\prime}(x)=\mathbf{K}^{\prime}(x) u(x),
$$

where $f$ and $u$ are defined in Lemma 3. Note that $f_{1}$ strictly decreases in $0<$ $t<1$. Indeed,

$$
f_{1}^{\prime}(x)=-\frac{2 \mathbf{E}^{\prime 2}(x) f_{2}(x)}{x\left(1-x^{2}\right)}
$$

where

$$
f_{2}(x)=x\left(1+x^{2}\right)\left(\frac{\mathbf{K}^{\prime}(x)}{\mathbf{E}^{\prime}(x)}\right)^{2}+(x-1)^{2} \frac{\mathbf{K}^{\prime}(x)}{\mathbf{E}^{\prime}(x)}-2
$$

Since $\mathbf{K}^{\prime}(x) / \mathbf{E}^{\prime}(x)>2 /\left(1+x^{2}\right)$, by (3.12), it follows that

$$
f_{2}(x)>\frac{4 x}{1+x^{2}}+\frac{2(1-x)^{2}}{1+x^{2}}-2=0 .
$$

Therefore, $f_{1}^{\prime}(x)<0$ for $0<x<1$. Hence, $f_{1}$ is strictly decreasing on this interval.
Finding the limit values of $f_{1}(x)$, we obtain

$$
\lim _{x \rightarrow+0} f_{1}(x)=+\infty, \quad \lim _{x \rightarrow 1-0} f_{1}(x)=-\pi^{2}
$$

Lemma 3 implies that $f_{1}\left(x_{1}\right) f_{1}\left(x_{2}\right)<0$ for all $t, 0<t<t_{0}$. Therefore, $G^{\prime}(t)<0$ if and only if

$$
\begin{equation*}
f_{1}\left(x_{1}(t)\right)+f_{1}\left(x_{2}(t)\right)>0 . \tag{3.17}
\end{equation*}
$$

To prove that (3.17) holds for $0<t \leq t_{*}$, we consider iterations $t_{n}$, where $t_{1}=0$ and $t_{n+1}$ is defined inductively by the following chain:

$$
\begin{aligned}
t_{n} & \mapsto x_{2, n}=x_{2}\left(t_{n}\right) \\
& \mapsto f_{1, n}=-f_{1}\left(x_{2, n}\right) \\
& \mapsto x_{1, n}=f_{1}^{-1}\left(f_{1, n}\right) \\
& \mapsto t_{n+1}=x_{1}^{-1}\left(x_{1, n}\right) .
\end{aligned}
$$

The results of 12 iterations are shown in Table 1.

| $n$ | 1 | 2 | 3 | 4 | 5 | 6 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $t_{n}$ | 0 | 3.64142 | 3.90656 | 3.98852 | 4.02433 | 4.04314 |
| $x_{2, n}$ | 1 | 0.14002 | 0.09653 | 0.08070 | 0.07258 | 0.06765 |
| $f_{1, n}$ | $\pi^{2}$ | 5.46528 | 3.78714 | 2.89890 | 2.34856 | 1.97399 |
| $x_{1, n}$ | 0.01174 | 0.02079 | 0.02638 | 0.03009 | 0.03271 | 0.03466 |
| $t_{n+1}$ | 3.64142 | 3.90656 | 3.98852 | 4.02433 | 4.04314 | 4.05424 |


| $n$ | 7 | 8 | 9 | 10 | 11 | 12 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $t_{n}$ | 4.05424 | 4.06133 | 4.06613 | 4.06954 | 4.07204 | 4.07393 |
| $x_{2, n}$ | 0.06435 | 0.06198 | 0.06020 | 0.05881 | 0.05770 | 0.05680 |
| $f_{1, n}$ | 1.70253 | 1.49674 | 1.33535 | 1.20539 | 1.09848 | 1.00900 |
| $x_{1, n}$ | 0.03616 | 0.03736 | 0.03833 | 0.03913 | 0.03981 | 0.04039 |
| $t_{n+1}$ | 4.06133 | 4.06613 | 4.06954 | 4.07204 | 4.07393 | 4.07539 |

Table 1. Iterations $t_{n}, x_{2, n}, f_{1, n}, x_{1, n}$.

We should emphasize that all explicitly shown digits of all entries in Table 1 are exact. The computed functions $x_{1}(t), x_{2}(t), f_{1}(x)$ and their inverses are real analytic. The known computational methods allow us to compute their values at a finite number of points (here 48) with the desired accuracy. Throughout the paper, we keep 5 exact digits, but the computations were actually done up to 10 digits. All computations were made with Maple and then doublechecked with Mathematica.
Note that

$$
t_{12}=4.07394 \ldots>t_{*}=4.07389 \ldots
$$

Let $t^{\prime}, 0<t^{\prime} \leq t_{*}$ satisfy $t_{n}<t^{\prime} \leq t_{n+1}$. Then, since the functions $x_{1}, x_{2}, f_{1}$ are monotone, we have:

$$
\begin{array}{rlrl}
x_{1, n-1}<x_{1}\left(t^{\prime}\right) & \leq x_{1, n}, & x_{2, n+1} \leq x_{2}\left(t^{\prime}\right) & <x_{2, n} \\
-f_{1, n}<f_{1}\left(x_{2}\left(t^{\prime}\right)\right) & \leq-f_{1, n+1}, & f_{1, n} \leq f_{1}\left(x_{1}\left(t^{\prime}\right)\right)<f_{1, n-1}
\end{array}
$$

for $n=1, \ldots, 16$, where $x_{1,0}=0$ and $f_{1,0}=+\infty$. Hence, $f_{1}\left(x_{1}\left(t^{\prime}\right)\right)+f_{1}\left(x_{2}\left(t^{\prime}\right)\right)>0$. Therefore, $G^{\prime}(t)<0$ and $G(t)$ strictly decreases in $0<t \leq t_{*}$.
Part 3: Now we show that for a fixed $r, 0<r<1$, and $a$ defined by (2.7), the limit point $b=a$ provides a local minimum for $\Omega(a, b)$. We put $b=a-\delta$ with $\delta>0$ small enough and then estimate the arctangents in (2.8). The well known upper bound

$$
\mu(s)<\log \frac{4}{s}
$$

see [AVV, p. 80] leads to the following estimate for $\tau_{1}=\tau_{1}(a, a-\delta)$ :

$$
0<\tau_{1}<4 e^{-\pi / \delta}
$$

which shows that

$$
0 \leq \tan ^{-1} \sqrt{\tau_{1}(a, a-\delta)}<2 e^{-\pi /(2 \delta)}=o\left(\delta^{n}\right)
$$

for any positive integer $n$ as $\delta \rightarrow+0$.
Since the second arctangent in (2.8) is differentiable at $b=a$, we obtain

$$
\begin{equation*}
\Omega(a, a-\delta)=\Omega(a, a)-\frac{\delta \mu^{2}\left(\tau_{2}\right)}{\pi^{2} \mu^{\prime}\left(\tau_{2}\right)} \frac{1}{\tau_{2}^{1 / 2}\left(1+\tau_{2}\right)}+o(\delta) \tag{3.18}
\end{equation*}
$$

as $\delta \rightarrow+0$. Since $\mu^{\prime}(\tau)<0$ for $0<\tau<1$, (3.18) shows that the degenerate configuration provides a local minimum in the problem under consideration.
Part 4: Since $G(t)$ and $\mu\left(r^{2}\right)$ both are monotone, it follows that (3.13) defines a decreasing function $t=t(r)$ from $r_{*} \leq r<1$ onto $0<t \leq t_{*}$, where $r_{*}=$ $.21886 \ldots$ is a solution of the equation

$$
\frac{2}{\mu\left(r^{2}\right)}=\frac{1}{\mu\left(x_{1}\left(t_{*}\right)\right)}+\frac{1}{\mu\left(x_{2}\left(t_{*}\right)\right)} .
$$

Assume now that for some $r^{\prime}, 0<r^{\prime}<1$ there is a nonsymmetric nondegenerate extremal configuration. Let $t^{\prime}, 0<t^{\prime}<t_{0}$ be a corresponding value of $t$. Since we have shown in Part 1 that for $\hat{t} \leq t \leq t_{0}$ there is no non-symmetric nondegenerate configuration, it follows that $0<t^{\prime}<\hat{t}$.
Note that

$$
\begin{equation*}
G(t) \leq G(\hat{t}) \quad \text { for } \hat{t} \leq t \leq t_{0} \tag{3.19}
\end{equation*}
$$

Indeed, $G(t)$ strictly decreases in $\hat{t}<t \leq t_{*}$ and

$$
G(t) \leq \frac{1}{\mu\left(x_{0}\right)}+\frac{1}{\mu\left(x_{2}\left(t_{*}\right)\right)}=G(\hat{t})
$$

for $t_{*} \leq t \leq t_{0}$.
Since $G\left(t^{\prime}\right)=2 / \mu\left(r^{\prime 2}\right)>G(\hat{t})$, (3.19) and the monotonicity property of $G(t)$ for $0<t \leq t_{*}$ imply that for fixed $r=r^{\prime}$, equation (3.13) has a unique solution $t=t^{\prime}$. Therefore, for $r=r^{\prime}$ there is only one critical point $b=b^{\prime}, 0<b^{\prime}<a$ of the constrained minimization problem (2.7), (2.8).

Since we have shown in Part 3 that $b=a$ provides a local minimum, it follows that $b=b^{\prime}$ provides a local maximum for the average harmonic measure considered as a function of $b$. Hence the symmetric and degenerate configurations are the only possible ones minimizing the average harmonic measure. Now the required assertion follows from Lemma 2.

The idea used to show that $G^{\prime}(t)<0$ for $0<t \leq t_{*}$ can be applied to prove monotonicity of $G$ for the whole range $0<t<t_{0}$. For this we first estimate $f_{1}\left(x_{1}(t)\right)+f_{1}\left(x_{2}(t)\right)$ using asymptotic expansions of $x_{1}(t)$ and $x_{2}(t)$ in a vicinity of $t=t_{0}$, then apply iterations described above. However, this requires more iterations and is not necessary for the proof of Theorem 1.
Some remarks concerning the symmetric case might be interesting. As we noticed in the begining of this section, $b=0$ is a critical point of the minimization problem under consideration. To decide whether it provides a minimum or maximum, we apply the standard second derivative test. From (2.7), (2.8), we find

$$
\begin{equation*}
\pi \frac{\partial^{2} \Omega}{\partial b^{2}}=\sum_{i=1}^{2}\left(\frac{1}{2} \frac{1+3 \tau_{i}}{\tau_{i}^{3 / 2}\left(1+\tau_{i}\right)^{2}}\left(\frac{\partial \tau_{i}}{\partial b}\right)^{2}-\frac{1}{\tau_{i}^{1 / 2}\left(1+\tau_{i}\right)} \frac{\partial^{2} \tau_{i}}{\partial b^{2}}\right) \tag{3.20}
\end{equation*}
$$

$$
\begin{align*}
\frac{\partial \tau_{i}}{\partial b} & =\frac{(-1)^{i+1} \mu^{2}\left(\tau_{i}\right)}{\pi \mu^{\prime}\left(\tau_{i}\right)} \\
\frac{\partial^{2} \tau_{i}}{\partial b^{2}} & =-\frac{\mu^{\prime \prime}\left(\tau_{i}\right) \mu\left(\tau_{i}\right)-2 \mu^{\prime 2}\left(\tau_{i}\right)}{\mu^{\prime}\left(\tau_{i}\right) \mu\left(\tau_{i}\right)}\left(\frac{\partial \tau_{i}}{\partial b}\right)^{2} \tag{3.21}
\end{align*}
$$

Let $\tau=\tau_{1}(a, 0)=\tau_{2}(a, 0)=r^{2}$. Then (3.20), (3.21) lead to

$$
\begin{align*}
\left.\pi \frac{\partial^{2} \Omega}{\partial b^{2}}\right|_{b=0}= & \frac{\mu^{3}(\tau)}{\pi^{2} \tau^{3 / 2}(1+\tau)^{2} \mu^{\prime 3}(\tau)}  \tag{3.22}\\
& \times\left[(1+3 \tau) \mu^{\prime}(\tau) \mu(\tau)+2 \tau(1+\tau)\left(\mu^{\prime \prime}(\tau) \mu(\tau)-2 \mu^{\prime 2}(\tau)\right)\right]
\end{align*}
$$

Using the differentiation formulas for $\mu(\tau)$, see [AVV, p. 82], and Legendre's relation

$$
\mathbf{E}(\tau) \mathbf{K}^{\prime}(\tau)+\mathbf{E}^{\prime}(\tau) \mathbf{K}(\tau)-\mathbf{K}(\tau) \mathbf{K}^{\prime}(\tau)=\frac{\pi}{2}
$$

we obtain from (3.22),

$$
\left.\pi \frac{\partial^{2} \Omega}{\partial b^{2}}\right|_{b=0}=\frac{1}{\pi^{2}} \tau^{1 / 2}(1-\tau) \mathbf{K}^{\prime 3}(\tau) u(\tau)
$$

with $u(x)$ defined by (3.9). Since $u(x)$ has only one zero $x=x_{0}$ in $0<x<1$, see Lemma 3, it follows that

$$
\begin{aligned}
& \left.\frac{\partial^{2} \Omega}{\partial b^{2}}\right|_{b=0}<0 \quad \text { if } \quad 0<r=\sqrt{\tau}<r_{0} \\
& \left.\frac{\partial^{2} \Omega}{\partial b^{2}}\right|_{b=0}>0 \quad \text { if } \quad r_{0}<r=\sqrt{\tau}<1
\end{aligned}
$$

which shows that for $0<r<r_{0}$ the symmetric configuration provides a local maximum and for $r_{0}<r<1$ it provides a local minimum. Figure 7 demonstrates typical behavior of the average harmonic measure $\Omega(a, b)$ as a function of $b$ for some values of $a$ or, equivalently, for some values of the main parameter $r$.


Figure 7. Typical behavior of $\Omega(a, b)$.

## 4. Some remarks and questions

Lemma 1, which actually describes the range of all possible pairs of the harmonic measures $\omega\left(r, E, D_{1}\right), \omega\left(-r, E, D_{2}\right)$, can be used to extremize some other characteristics, besides the average harmonic measure, of the studied configurations. In particular, a solution to the problem

$$
\begin{equation*}
\min _{E} \max \left\{\omega\left(r, E, D_{1}\right), \omega\left(-r, E, D_{2}\right)\right\} \tag{4.1}
\end{equation*}
$$

is an immediate consequence of Lemma 1 and the easy observation that the equality $\omega\left(r, E, D_{1}\right)=\omega\left(-r, E, D_{2}\right)$ must hold for any extremal configuration. This implies that the symmetric configuration is the unique minimizer of (4.1).
It is interesting to note that the $n$-dimensional counterparts of Problems (2.1) and (4.1) are much harder. They are open even in the case $n=3$, where both questions sound very natural. The difficulty here is the lack of an analogous theory to Jenkins' theory of the extremal decompositions into nonoverlapping regions, see [J1, S2].

A symmetrization result of A . Baernstein $[\mathrm{B}, \mathrm{BT}]$ that is still applicable, shows that extremal regions $D_{1}, D_{2}$ are bounded by surfaces of revolution about the axis containing both reference points $a_{1}=r$ and $a_{2}=-r$. Of course, for the extremal regions $D_{1}, D_{2}$ of the $n$-dimensional version of problem (4.1) the equality $\omega\left(a_{1}, E, D_{1}\right)=\omega\left(a_{2}, E, D_{2}\right)$ is satisfied in all dimensions.
Another interesting open question concerns a similar control for the average or maximal harmonic measure at $n \geq 3$ points $a_{k}=r e^{2 \pi i(k-1) / n}, k=1, \ldots, n$, equally distributed on the circle $\{z:|z|=r\}, 0<r<1$, in the unit disk $\Delta$. More precisely, the following problem would be of some interest.
Let a continuum $E$ split $\Delta$ into $n$ simply connected (or more general) regions $D_{k} \ni a_{k}, k=1, \ldots, n$. Find

$$
\inf _{E} \frac{1}{n} \sum \omega\left(a_{k}, E, D_{k}\right) \quad \text { and } \quad \inf _{E} \max _{k=1, \ldots, n} \omega\left(a_{k}, E, D_{k}\right)
$$

and describe possible extremal configurations.
The conjectured extremal configuration of the minmax problem consists of $n$ circular sectors

$$
S_{k}=\left\{z:|z|<1,\left|\arg z-2 \pi \frac{k-1}{n}\right|<\frac{\pi}{n}\right\}, \quad k=1, \ldots, n .
$$

Under an additional assumption that the closure of $\partial D_{k} \cap \mathbb{U}$ is connected for all $k=1, \ldots, n$, this can be proved applying Jenkins's theory on the extremal decomposition [J1, S2]. This result will be published elsewhere.

## References

A1. L. V. Ahlfors, Conformal Invariants, Topics in geometric function theory, McGraw-Hill Book Co., 1973.
A2. Lectures on Quasiconformal Mappings, Reprint of the 1966 original, Wadsworth \& Brooks/Cole Advanced Books \& Software, Monterey, CA, 1987.
AVV. G. D. Anderson, M. K. Vamanamurthy, and M. K. Vuorinen, Conformal Invariants, Inequalities, and Quasiconformal Maps, John Wiley \& Sons, Inc., New York, 1997.
B. A. Baernstein II, Integral means, univalent functions and circular symmetrization, Acta Math, 133 (1974), 139-169.
BT. A. Baernstein II and B. A. Taylor, Spherical rearrangements, subharmonic functions, and $*$-functions in $n$-space, Duke Math. J. 43 (1976), 245-268.
D. V. N. Dubinin, Symmetrization in geometric theory of functions of a complex variable, Uspehi Mat. Nauk 49 (1994), 3-76; English translation in Russian Math. Surveys 49 (1994) 1, 1-79.

Du. P. L. Duren, Univalent Functions, Springer-Verlag New York Inc., 1983.
G. G. M. Goluzin: Geometric Theory of Functions of a Complex Variable, AMS, Providence, R.I., 1969.
H. W. K. Hayman, Multivalent Functions, Cambridge Univ. Press, Cambridge, 1958.

He. J. Hersch, Longeurs extremales, mesure harmonique et distance hyperbolique (in French), C.R. Acad. Sci. Paris 235 (1952), 569-571.
J1. J. A. Jenkins, On the existence of certain general extremal metrics, Ann. of Math. (2) 66 (1957), 440-453.

J2. , Univalent Functions and Conformal Mappings, Springer-Verlag, 1958.
J3. , The method of the extremal metric, in: The Bieberbach conjecture (West Lafayette, Ind., 1985), Math. Surveys Monogr., 21, Amer. Math. Soc., Providence, RI, 1986, 95-104.
Kr. J. Krzyz, Circular symmetrization and Green's function, Bull. Acad. Polon. Sci. Sér. Sci. Math. Astr. Phys. 7 (1959), 327-330.
P. Ch. Pommerenke, Boundary Behaviour of Conformal Maps, Springer-Verlag, 1992.

S1. A. Yu. Solynin, Extremal problems on conformal moduli and estimates for harmonic measures, J. Analyse Math. 74 (1998), 1-49.
S2. $\qquad$ , Modules and extremal metric problems, Algebra i Analiz 11 (1999) no. 1, 3-86; English translation in St. Petersburg Math. J. 11 (2000), 1-65.

Roger W. Barnard
E-MAIL: barnard@math.ttu.edu Address: Department of Mathematics and Statistics, Texas Tech University, Lubbock, Texas 79409-1042, U.S.A.

Leah Cole
E-MAIL: colel@wbu.edu
Address: Wayland Baptist University, Division of Mathematics and Sciences, Plainview, TX 79072, U.S.A.

Alexander Yu. Solynin E-maIL: solynin@pdmi.ras.ru
Address: Steklov Institute of Mathematics at St. Petersburg, Russian Academy of Sciences, Fontanka 27, St. Petersburg, 191011, Russia

